

## **II. Ermakov's direct Method**

Objekttyp: **Chapter**

Zeitschrift: **L'Enseignement Mathématique**

Band (Jahr): **11 (1965)**

Heft 2-3: **L'ENSEIGNEMENT MATHÉMATIQUE**

PDF erstellt am: **12.07.2024**

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of the problem and prove generalized versions of Pringsheim's results.

This note brings therefore an improvement and simplification of the sections I-V and XI of [5], while I have nothing to add to the sections VI-X of [5].

## II. ERMAKOF'S DIRECT METHOD

6. The form of the expression (1) makes it plausible that we will have to use the integral transformation formula

$$\int_a^b f(\Psi(x)) \Psi'(x) dx = \int_{\Psi(a)}^{\Psi(b)} f(x) dx. \quad (3)$$

In order to be able to use (3) we have in any case to assume that  $f(x)$  is integrable in the integration interval and  $\Psi(x)$  totally continuous between  $a$  and  $b$ . However, additional conditions are necessary and two such conditions are known either of which ensures the relation (3):

- $J_1$ :  $|f(x)|$  is uniformly bounded in the integration interval;
- $J_2$ :  $\Psi(x)$  is monotonically increasing or monotonically decreasing.

7. THEOREM 1. Assume that  $\psi(x)$  and  $\Psi(x)$  are totally continuous for  $x \geq x_0$  and that we have for a sequence  $b_v \geq x_0$  ( $v = 1, 2, \dots$ )

$$\psi(b_v) \leq \Psi(b_v), \quad \Psi(b_v) \rightarrow \infty (v \rightarrow \infty). \quad (4)$$

Let  $f(x)$  be  $\geq 0$  on no half-line  $x \geq \xi$  almost everywhere  $= 0$ , and measurable in an interval  $J$  containing all values of  $\psi(x)$  and  $\Psi(x)$  for  $x \geq x_0$ . Assume further that for any finite subinterval of  $J$  the transformation formula (3) holds as well for  $\psi(x)$  as for  $\Psi(x)$ . Then, if we have for almost all  $x$  with  $x \geq x_0$  and for an  $\alpha$  with  $0 < \alpha < 1$ :

$$f(\Psi(x)) \Psi'(x) \leq \alpha f(\psi(x)) \psi'(x) (x \geq x_0), \quad 0 < \alpha < 1, \quad (5)$$

the integral (2) is convergent and we have for all  $x \geq x_0$ :

$$\Psi(x) > \psi(x) \quad (x \geq x_0). \quad (6)$$

8. *Proof.* For an arbitrary  $x \geq x_0$  integrate (5) between  $x$  and  $b_v > x$ . Then we have, using (3):

$$\int_{\Psi(x)}^{\Psi(b_v)} f(x) dx \leq \alpha \int_{\psi(x)}^{\psi(b_v)} f(x) dx$$

and this remains true, by (4), if  $\psi(b_v)$  is replaced by  $\Psi(b_v)$ . We can therefore write

$$\int_{\Psi(x)}^{\Psi(b_v)} f(x) dx \leq \alpha \int_{\Psi(x)}^{\Psi(b_v)} f(x) dx + \alpha \int_{\psi(x)}^{\Psi(x)} f(x) dx,$$

or, bringing the first right hand term to the left:

$$(1 - \alpha) \int_{\Psi(x)}^{\Psi(b_v)} f(x) dx \leq \alpha \int_{\psi(x)}^{\Psi(x)} f(x) dx.$$

But here, if we take  $x = b_1$  it follows for  $b_v \rightarrow \infty$  the convergence of (2) and also that the right hand expression is  $> 0$  for any  $x \geq x_0$ . (6) follows immediately and the Theorem 1 is proved.

9. THEOREM 2. Assume that  $\psi(x)$ ,  $\Psi(x)$  are totally continuous for  $x \geq x_0$  and that  $f(x)$  is non-negative and measurable in an interval  $J$  containing all values of  $\psi(x)$  and  $\Psi(x)$ . Assume that (3) holds as well for  $\psi(x)$  as for  $\Psi(x)$ . Assume further that there exists an  $a \geq x_0$  such that:

$$\int_{\psi(a)}^{\Psi(a)} f(x) dx > 0, \quad (7)$$

and a sequence  $b_v \geq x_0$  ( $v = 1, 2, \dots$ ) such that:

$$\psi(b_v) \rightarrow \infty, \quad \Psi(b_v) \rightarrow \infty (v \rightarrow \infty). \quad (8)$$

Then if we have for almost all  $x \geq x_0$ :

$$f(\Psi(x)) \Psi'(x) \geq f(\psi(x)) \psi'(x), \quad (9)$$

the integral (2) is divergent and we have for all  $x \geq a$ :

$$\Psi(x) > \psi(x) \quad (x \geq a). \quad (10)$$

10. *Proof.* For any  $x > a$  we obtain from (9), integrating on both sides from  $a$  to  $x$  and using (3):

$$\int_{\Psi(a)}^{\Psi(x)} f(x) dx \geq \int_{\psi(a)}^{\psi(x)} f(x) dx$$

and therefore

$$\int_{\psi(x)}^{\Psi(x)} f(x) dx \geq \int_{\psi(a)}^{\Psi(a)} f(x) dx \quad (x \geq a). \quad (11)$$

This proves already (10).

Putting in (11)  $x = b_v$  it follows

$$\int_{\psi(b_v)}^{\Psi(b_v)} f(x) dx \geq \int_{\psi(a)}^{\Psi(a)} f(x) dx \quad (12)$$

while, if (2) were convergent, the left side integral in (12) would tend to 0.

Theorem 2 is proved.

11. THEOREM 3. Assume that  $\psi(x)$  and  $\Psi(x)$  are totally continuous for  $x \geq x_0$ , that (3) holds as well for  $\psi(x)$  as for  $\Psi(x)$  and  $f(x)$  is  $\geq 0$  and measurable in an interval containing all values of  $\psi(x)$  and  $\Psi(x)$  for  $x > x_0$  without being almost everywhere = 0 in  $(\Psi(a), \infty)$ . Assume further that there exists a constant  $\gamma$ ,  $0 < \gamma < 1$ , and a sequence  $b_v \geq x_0$  ( $v = 1, 2, \dots$ ) such that

$$\gamma\psi(b_v) \leq \Psi(b_v), \quad 0 < \gamma < 1, \quad \psi(b_v) \rightarrow \infty (v \rightarrow \infty), \quad (13)$$

and further that for a constant  $c$  from a certain  $x = x_1 \geq x_0$  on:

$$f(x) \leq \frac{c}{x} \quad (x \geq x_1). \quad (14)$$

Assume finally that for a constant  $\alpha$ ,  $0 < \alpha < 1$ :

$$f(\Psi(x)) \Psi'(x) \leq \alpha f(\psi(x)) \psi'(x), \quad 0 < \alpha < 1. \quad (15)$$

Then the integral (2) converges and we have  $\Psi(x) > \psi(x)$  for all  $x > x_0$ .

12. *Proof.* We have as in the proof of the Theorem 1:

$$\int_{\Psi(x_0)}^{\Psi(b_v)} f(x) dx \leq \alpha \int_{\psi(x_0)}^{\psi(b_v)} f(x) dx$$

and therefore, using (13)

$$\frac{\gamma\psi(b_v)}{\Psi(x_0)} \int_{\Psi(x_0)}^{\psi(b_v)} f(x) dx \leq \alpha \int_{\psi(x_0)}^{\psi(b_v)} f(x) dx = \alpha \int_{\psi(x_0)}^{\psi(b_v)} f(x) dx + \alpha \int_{\psi(x_0)}^{\psi(b_v)} f(x) dx.$$

But the last right hand integral is, by (14),  $\leq c \log \frac{1}{\gamma}$ , so

that we obtain:

$$(1 - \alpha) \int_{\Psi(x_0)}^{\psi(b_v)} f(x) dx \leq \int_{\psi(x_0)}^{\Psi(x_0)} f(x) dx + c \log \frac{1}{\gamma}.$$

The convergence of (2) follows now immediately from  $\psi(b_v) \rightarrow \infty$ .

13. Suppose that we have, on the other hand, for an  $a > x_0$ :

$$\Psi(a) \leq \psi(a).$$

Proceeding then as in the proof of the Theorem 1 we have, as from  $\psi(b_v) \rightarrow \infty$  and the total continuity of  $\psi(x)$  follows  $b_v \rightarrow \infty$ , for  $b_v \geq a$ :

$$\int_{\Psi(a)}^{\Psi(b_v)} f(x) dx \leq \alpha \int_{\psi(a)}^{\psi(b_v)} f(x) dx,$$

and, for  $v \rightarrow \infty$ :

$$\int_{\Psi(a)}^{\infty} f(x) dx \leq \alpha \int_{\psi(a)}^{\infty} f(x) dx.$$

But here the left hand integral is  $> 0$ , the right hand integral is majorized by it and the relation is impossible for  $\alpha < 1$ .<sup>3)</sup>

### III. A NEW METHOD FOR NOT NECESSARILY MONOTONIC $f(x)$

14. THEOREM 4. Assume that  $\Psi(x)$  is for  $x \geq x_0$  a positive and monotonically increasing differentiable function for which

<sup>3)</sup> Observe that in Ermakov's paper [1] the criteria are given in the following form:  
 $\sum_{v=1}^{\infty} f(v)$  for a monotonic  $f(x)$  is convergent or divergent according as

$$\lim_{x \rightarrow \infty} \frac{f(\Psi(x))\Psi'(x)}{f(\psi(x))\psi'(x)}$$

is  $< 1$  or  $> 1$ . In the note [2] Ermakov takes  $\Psi(x) \equiv x$  which is no essential specialisation. However, the conditions (5) for convergence and (9) for divergence (with the specialisation  $\Psi(x) \equiv x$ ) are already found in the textbooks, see e.g. [3].