Particles exist in the low temperature ^4_2 model

Autor(en): Koch, H.

Objekttyp: Article

Zeitschrift: Helvetica Physica Acta

Band (Jahr): 53 (1980)

Heft 3

PDF erstellt am: 08.08.2024

Persistenter Link: https://doi.org/10.5169/seals-115129

Nutzungsbedingungen

Die ETH-Bibliothek ist Anbieterin der digitalisierten Zeitschriften. Sie besitzt keine Urheberrechte an den Inhalten der Zeitschriften. Die Rechte liegen in der Regel bei den Herausgebern. Die auf der Plattform e-periodica veröffentlichten Dokumente stehen für nicht-kommerzielle Zwecke in Lehre und Forschung sowie für die private Nutzung frei zur Verfügung. Einzelne Dateien oder Ausdrucke aus diesem Angebot können zusammen mit diesen Nutzungsbedingungen und den korrekten Herkunftsbezeichnungen weitergegeben werden.

Das Veröffentlichen von Bildern in Print- und Online-Publikationen ist nur mit vorheriger Genehmigung der Rechteinhaber erlaubt. Die systematische Speicherung von Teilen des elektronischen Angebots auf anderen Servern bedarf ebenfalls des schriftlichen Einverständnisses der Rechteinhaber.

Haftungsausschluss

Alle Angaben erfolgen ohne Gewähr für Vollständigkeit oder Richtigkeit. Es wird keine Haftung übernommen für Schäden durch die Verwendung von Informationen aus diesem Online-Angebot oder durch das Fehlen von Informationen. Dies gilt auch für Inhalte Dritter, die über dieses Angebot zugänglich sind.

Ein Dienst der *ETH-Bibliothek* ETH Zürich, Rämistrasse 101, 8092 Zürich, Schweiz, www.library.ethz.ch

http://www.e-periodica.ch

Helvetica Physica Acta, Vol. 53 (1980), Birkhäuser Verlag, Basel

Particles exist in the low temperature φ_2^4 model

by H. Koch

Département de Physique Théorique, Université de Genève, 1211 Genève 4, Switzerland

(22. VIII. 1980)

Abstract. The existence of an upper mass gap and of a two particle bound state are shown for the $\lambda \varphi^4 - \frac{1}{4}\varphi^2$ field theory in two dimensions. The proof is based on a resummation of the low temperature cluster expansion which can be combined with Spencer's expansion in order to obtain decay properties of irreducible kernels.

I. Introduction and main results

Once the existence of some quantum field model is established (in the sense of [OSch]), it is natural to analyze physically relevant quantities such as the mass spectrum, the S-matrix etc. This has for example been done for $\lambda \Re(\varphi)_2 + \frac{1}{2}m_0^2\varphi^2$ models for small coupling λ [S], [SZ], [DE], [OS], [K], [GJ] after Glimm, Jaffe and Spencer [GJS I, II] had constructed these models by introducing the cluster expansion.

In the meantime much progress has been made in understanding field theories at low temperatures λ^{-1} . For the two dimensional $\lambda \varphi^4 - \frac{1}{4}\varphi^2 - \mu \varphi$ models with $|\mu| < \lambda^2 \ll \lambda$ (which is for $\mu = 0$ equivalent to a $\lambda' \varphi^4 + \frac{1}{2} \varphi^2$ model with $\lambda' \gg 1$), Glimm, Jaffe and Spencer developed a convergent expansion which establishes existence and mass gap of each pure phase associated to a minimum of the polynomial. Their methods have also been applied to other low temperature models [Br], [BF], [BG], [Su].

In this paper we study the mass spectrum in a pure phase of the Euclidean $\Re(\varphi) = \lambda \varphi^4 - \frac{1}{4}\varphi + (64\lambda)^{-1}$ model with $0 \le \lambda \ll 1$. Our method is a minimal synthesis of methods used in the one phase region [S], [K], with low temperature expansions [GJS III], [BG]. This involves a proof of n+1 particle decay for n particle irreducible kernels, n = 0, 1, 2. The upper gap then follows by a result of Burnap [B] which is model independent.

In the two pure phases the mean of the field φ takes a value near $\pm \xi$, where \mathfrak{P} has its minimum, $\xi = (8\lambda)^{-1/2}$. In both cases the interaction is $\int_{\Lambda} d^2 x : \mathfrak{P}(\varphi) : (x)$ in the limit $\Lambda \to \mathbb{R}^2$. We will choose the + phase by taking a large region $Y \supset \Lambda$ and imposing suitable boundary conditions on the field in the strip $Y \setminus \Lambda$.

The polynomial \mathfrak{P} can be written in terms of the translated field $\phi = \varphi - \xi$,

$$\mathfrak{P}(\varphi) = \lambda \phi^4 + (2\lambda)^{1/2} \phi^3 + \frac{1}{2} \phi^2 = V(\phi) + \frac{1}{2} \phi^2.$$

Let

$$\mathbf{V}(\Lambda) = \int_{\Lambda} d^2 \mathbf{x} : \mathbf{V}(\boldsymbol{\phi}) : (\mathbf{x})$$

where : : denotes Wick ordering with respect to the covariance $C = (-\Delta_{\partial Y} + 1)^{-1}$ with $\Delta_{\partial Y}$ being the Laplacean with zero Dirichlet data on ∂Y . Then we define the $\lambda \varphi^4 - \frac{1}{4}\varphi^2 - (64\lambda)^{-1}$ quantum field theory by the family of Schwinger functions

$$S(x_1, \dots, x_n) = \langle \phi(x_1) \cdots \phi(x_n) \rangle = \frac{\int d\phi_C \phi(x_1) \cdots \phi(x_n) e^{-v(\Lambda)}}{\int d\phi_C e^{-v(\Lambda)}}$$

in the limit $\Lambda \to \mathbb{R}^2$. As usually $d\phi_C$ denotes the Gaussian process with mean zero and covariance C. A condition on the sequence Λ , Y is that $Y \supset N(\Lambda)$, where for $X \subset \mathbb{R}^2$

$$N(X) = \{x \in \mathbb{R}^2 : dist (x, X) \le L\}$$

and $L \in \mathbb{N}$, $L \simeq (\log \lambda)^2$ is a given length scale.

Our main result is as follows.

Theorem 1. Fix $\varepsilon > 0$. Then there is a $\delta > 0$ such that the following is true.

- (a) The infinite volume Schwinger functions exist and are analytic in λ for $|\text{Im }\lambda| < \delta \text{ Re }\lambda < \delta^3$.
- (b) They are also C^{∞} in $\lambda^{1/2}$ for $\lambda^{1/2} \in [0, \delta]$ and define a Wightman field theory.
- (c) For $\lambda^{1/2} \in [0, \delta]$ the mass spectrum in the interval $[0, 2(1-\varepsilon)]$ consists of two points, 0 and m, where $|m-1| < \varepsilon$, corresponding to the vacuum and the one particle state respectively. Furthermore $m = m(\lambda^{1/2})$ is C^{∞} in $\lambda^{1/2}$ and

$$m^{2}(\lambda^{1/2}) = 1 - 4\sqrt{3} \lambda + O(\lambda^{2}).$$

(d) If $\lambda^{1/2} \in [0, \delta]$ then there is a two particle bound state with mass $m_B \in (2(1-\varepsilon), 2m)$. Furthermore $m_B = m_B(\lambda^{1/2})$ is C^{∞} in $\lambda^{1/2}$ and

$$m_{\rm B}^2(\lambda^{1/2}) = 4m^2(\lambda^{1/2}) - 144\lambda^2 + O(\lambda^{5/2}).$$

Let P be the spectral projection of the mass operator associated to a subinterval of (0, 2m) in the complement of $\{m, m_B\}$. Then $P\varphi(f) \times (1+\varphi(g))\Omega = 0$ for any pair f, g of C^{∞} functions with compact support.

The first part (a) and (b) of this theorem has already been proven by [GJS III], even for small external field. Our analysis differs on this level only by the way of dividing unnormalized expectations by the partition function.

Differentiability with respect to $\lambda^{1/2}$ can be shown as usually by introducing cutoffs. The explicit formula is

$$\partial_{\lambda^{1/2}} \langle \Phi \rangle = -\int d^2 x [\langle \Phi \cdot (2\lambda^{1/2} : \phi^4 : (\mathbf{x}) + \sqrt{2} : \phi^3 : (\mathbf{x})) \rangle \\ - \langle \Phi \rangle \cdot \langle 2\lambda^{1/2} : \phi^4 : (\mathbf{x}) + \sqrt{2} : \phi^3 : (\mathbf{x}) \rangle].$$
(I.1)

The proof of (c) and (d) is based on decay properties of the following irreducible kernels. Denote by $S_{k,l}^1$ the integral operator defined by the kernel

$$S_{k,l}(x_1, \ldots, x_k; y_1, \ldots, y_l) = S(x_1, \ldots, y_l) - S(x_1, \ldots, x_k)S(y_1, \ldots, y_l)$$

and denote by $R_{21}(x_1, x_2; y)$ the truncated three point function. The n particle irreducible (n p.i.) kernels $S_{k,l}^{n+1}$ are inductively defined by

$$\mathbf{S}_{k,l}^{n+1} = \mathbf{S}_{k,l}^{n} (\mathbf{S}_{n,n}^{n})^{-1} \mathbf{S}_{n,l}^{n}.$$
(I.2)

Three other kernels are needed for our analysis. Namely the 1 p.i. two point function k, the 2 p.i. three point function L and the Bethe-Salpeter kernel K. They are defined by the equations

$$S_{1,1}^{1} = C - CkS_{1,1}^{1},$$

$$R_{21} = S_{2,2}^{2}LS_{1,1}^{1},$$

$$S_{2,2}^{2} = 2S_{1,1}^{1} \otimes S_{1,1}^{1} - 2(S_{1,1}^{1} \otimes S_{1,1}^{1})KS_{2,2}^{2}.$$

(I.3)

We intend to prove the following proposition.

Proposition 2. Let p > 1 and $\varepsilon > 0$ be given. Then for $\lambda > 0$ sufficiently small,

(a)
$$S_{k,l}^{n}(x_1, \ldots, x_k; y_1, \ldots, y_l) = O_p(\exp(-(1-\varepsilon)\min_{i,j} |x_i^0 - y_j^0|)), \quad n = 1, 2, 3$$

(b) $k(x; y) = O_{\infty}(\exp(-2(1-\varepsilon)|x^0-y^0|))$

(c) $L(x_1, x_2; y) = O_{\infty}(\exp(-(1-\varepsilon)[\frac{3}{2}|2y^0 - x_1^0 - x_2^0| + \frac{1}{2}|x_1^0 - x_2^0|]))$

(d)
$$K(x_1, x_2; y_1, y_2) = O_{\infty}(\exp(-(1-\varepsilon)[\frac{3}{2}|y_1^0 + y_2^0 - x_1^0 - x_2^0| + \frac{1}{2}|x_1^0 - x_2^0| + \frac{1}{2}|y_1^0 - y_2^0|]))$$

where $g(\mathbf{z}) = O_p(h(\mathbf{z}))$ means that $|\int d\mathbf{z}g(\mathbf{z})h(\mathbf{z})^{-1}| \le \text{const.} ||g||_{L_p}$ for every continuous function f with support in a product of unit squares $\Box_j = [j^0, j^0 + 1] \times [j^1, j^1 + 1]$.

These bounds can now be used to prove the second half of Theorem 1. By a result of Burnap [B] the decay of $S_{k,l}^n$ for n = 1, 2 and of K imply the upper mass gap. This, together with the remaining decay properties and differentiability in $\lambda^{1/2}$, allows to analyze the bound state spectrum exactly as in the $\lambda \Re(\varphi)_2$ case, $\lambda \ll 1$. However, the absence of other spectrum in $(2(1-\varepsilon), 2m)$ does not follow. This could probably be analyzed, as in the weak coupling region, by an n particle cluster expansion [GJS II]. It is replaced by the weaker statement in Theorem 1(d).

We shall now prepare the proof of Proposition 2. In a first step we combine the Peierls expansion with Spencer's t-expansion [S II]. Let $l \simeq |\log \lambda|^{1/4}$, $l \in \mathbb{N}$ such that $L \simeq (\log \lambda)^2$ is a multiple of *l*. By Δ we denote squares $l \times l$ with corners in $l\mathbb{Z}^2$, while unit squares with corners in \mathbb{Z}^2 will be denoted by \Box . We suppose Λ and Y to be simply connected unions of squares Δ .

In the field theoretic version of the Peierls expansion [GJS III] the mean field $\Box \mapsto \bar{\varphi}_{\Box} = \int_{\Box} d^2 x \varphi(x)$ plays the role of an Ising spin variable $\Box \mapsto \Sigma(\Box) = \pm$ associated to unit squares \Box in Λ . Let $1 = \chi_+ + \chi_-$ be the smooth partition of unity defined by

$$\chi_{+}(\mathbf{x}) = \chi_{-}(-\mathbf{x}) = \pi^{-1/2} \int_{0}^{\infty} d^{2} y e^{-(\mathbf{x}-\mathbf{y})^{2}}$$

Then we split the measure $e^{-v} d\varphi$ into $2^{|\Lambda|}$ parts by inserting

$$1 = \prod_{\Box \subset \Lambda} \left[\chi_{+}(\bar{\varphi}_{\Box}) + \chi_{-}(\bar{\varphi}_{\Box}) \right] = \sum_{\Sigma \in \{+,-\}^{|\Lambda|}} \chi_{\Sigma}$$
(I.4)

where

$$\chi_{\Sigma} = \prod_{\Box \subset \Lambda} \chi_{\Sigma(\Box)}(\bar{\varphi}_{\Box}).$$

Each term in the sum (I.4) represents a partition of Λ into pure phases, i.e. regions where Σ has a definite sign. By setting $\Sigma(\Box) = +$ for $\Box \subset Y \setminus \Lambda$ this sign is uniquely determined by the phase boundaries, which we will also denote by Σ .

Notice that it is very improbable for the field φ to take values near zero, or to change from $\pm \xi$ to $\pm \xi$. This will make the Σ expansion converge. On the other hand because of the factor χ_{Σ} the effective potential in a pure phase is near to $\frac{1}{2}(\varphi \pm \xi)^2$. This will be used to perform a cluster expansion in the region $Y \setminus N(\Sigma)$. For this purpose we translate the field $\phi = \varphi - \xi$ by a C[∞] function $\xi - g$ satisfying

 $g(x) = \Sigma(\Box) \cdot \xi$ for $x \in \Box$ with dist $(\Box, \Sigma) > \frac{1}{2}L$.

For the explicit form of g see [GJS III]. The expansion in phase boundaries for $\langle \Phi \rangle$ is then as follows.

$$\langle \Phi \rangle = \frac{\sum_{\Sigma} \int d\psi_{C} e^{-F(\Lambda,\Sigma)} \Phi \chi_{\Sigma} e^{-V(\Lambda)}}{\sum_{\Sigma} \int d\psi_{C} \chi_{\Sigma} e^{-V(\Lambda) - F(\Lambda,\Sigma)}}$$
(I.5)

where $\psi = \varphi - g$ and

$$F(\Lambda, \Sigma) = \log \frac{d\psi_{\rm C}}{d\phi_{\rm C}}$$

$$= \frac{1}{2} \langle (\xi - g), (-\Delta + 1)(\xi - g) \rangle + \langle \phi, (-\Delta + 1)(\xi - g) \rangle.$$
(I.6)

This is the starting point for the cluster expansion which will be defined later and which allows to take the limit $\Lambda \rightarrow \mathbb{R}^2$.

We shall now illustrate Spencer's t-expansion. It was designed to prove Proposition 2 in the one phase region. Let $\Delta_{\mathfrak{B}}$ be the Laplacean with zero Dirichlet data on a closed curve \mathfrak{L} in $\mathbb{R}^2 \cup \{\infty\}$, and for $t \in [0, 1]$ let

$$C(t) = (1-t)(-\Delta_{\mathfrak{L}}+1)^{-1} + t \cdot (-\Delta+1)^{-1}.$$

Notice that $C(t; x, y) = O(e^{-(1-\varepsilon)|x-y|})$, and C(t; x, y) = 0 if x and y are separated by \mathfrak{L} . Then for every $G_1, G_2 \subset \mathbb{R}^2$ separated by \mathfrak{L} and for each finite set A of pairs a = (x, y) with $x, y \in G_1 \cup G_2$ (for simplicity we assume $|x-y| \ge \text{const.}$) let us define $K_t = \prod_{a \in A} C(t; a)$. It is easy to see that

$$\partial_t^{\alpha} \mathbf{K}_t \Big|_{t=0} = 0 \quad \text{for} \quad \alpha < \mathbf{n}$$
 (I.7)

implies

$$|\mathbf{K}_{t}| \le O(1)^{|\mathbf{A}|} \exp(-n(1-\varepsilon) \operatorname{dist}(\mathbf{G}_{1},\mathbf{G}_{2})).$$
 (I.8)

But if we take A such that K_t is a graph of a $\mathfrak{P}(\varphi)_2$ theory, then $(I.7) \Rightarrow (I.8)$ is in general not true. In this case irreducibility should be checked at each line \mathfrak{P} separating G_1 from G_2 . However if the vertices produce small contributions, then it is sufficient to consider a certain discrete set of lines \mathfrak{P}_i , $i \in I$, depending on G_1 and G_2 . For $t = (t_i)_{i \in I}$ the corresponding covariance is defined by

$$C(t) = \prod_{i \in I} \left[(1 - t_i) \theta_{\mathfrak{L}_i} + t_i \right] (-\Delta_{\partial y} + 1)^{-1}$$

where $\theta_{\mathfrak{L}}(-\Delta_{\Gamma}+1)^{-1} = (\Delta_{\Gamma \cup \mathfrak{L}}+1)^{-1}$. t-derivatives are indexed by functions $\alpha: I \to \mathbb{Z}_+$,

$$\partial_t^{\alpha} \mathbf{K}_t = \prod_{i \in \mathbf{I}} \partial_{t_i}^{\alpha(i)} \mathbf{K}_t.$$

It was a goal of Spencer [S] to establish a relation analogous to $(I.7) \Rightarrow (I.8)$ for kernels K_t of a weakly coupled $\lambda \Re(\varphi)_2 + \frac{1}{2}m_0^2\varphi^2$ theory.

In trying to apply the same method to our problem, the following difficulty arises. In order to show (I.7) using an explicit formula for t derivatives and a factorization property (see (I.13), (I.14)), that t dependent covariance has to be introduced before doing the translation $\phi \rightarrow \psi$ (notice that g(x) depends on $\Sigma(\Box)$ if dist $(x, \Box) < \frac{1}{2}L$). But if we do so it is impossible to translate the field in the case where $\xi - g \neq 0$ on \mathfrak{L}_i and $t_i \neq 1$.

This problem can be solved by taking advantage of the small factors associated to nonempty intersections $\Sigma \cap N(\mathfrak{L}_i)$, so that we need not to consider $t_i \neq 1$. We perform a resummation of the Peierls expansion (r-expansion) by specifying phase boundary free region $(r_i = 0)$, where Spencer's expansion can be applied. In the complement of these pure phases, small factos (r-derivatives) associated to phase boundaries will compensate the missing convergence factors from t derivatives.

Let \mathfrak{F} be a finite subset of $l\mathbb{Z}$, to be determined later. To each $i \in \mathfrak{F}$ we associate the Dirichlet lines $\mathfrak{L}_i = \{(x^0, x^1) \in \mathbb{R}^2 : x^0 = i\}$ and an expansion in two variables

$$F(\mathbf{r}_{i} = 1, t_{i} = 1) = F(0, 0) + \int_{0}^{1} d\mathbf{r}_{i} \partial_{\mathbf{r}_{i}} F(\mathbf{r}_{i}, 1) + \int_{0}^{1} dt_{i} \partial_{t_{i}} F(0, t_{i})$$
(I.9)

applied to the expectation

6

$$\langle \Phi \rangle_{\mathbf{r},\mathbf{t}} = \frac{\sum_{i \in \mathfrak{Y}} \prod_{i \in \mathfrak{Y}} r_{i}^{|\Sigma \cap N(\mathfrak{Q}_{i})|} \int d\phi_{C(\mathbf{t})} \Phi \chi_{\Sigma} e^{-V(\Lambda)}}{\sum_{\Sigma} \prod_{i \in \mathfrak{Y}} r_{i}^{|\Sigma \cap N(\mathfrak{Q}_{i})|} \int d\phi_{C(\mathbf{t})} \chi_{\Sigma} e^{-V(\Lambda)}}$$
(I.10)

where $0^0 = 1$, $r = r^{\mathfrak{F}} = (r_i)_{i \in \mathfrak{F}}$ and $t = t^{\mathfrak{F}} = (t_i)_{i \in \mathfrak{F}}$. Notice that setting $r_i = 0$ eliminates all terms in the sum over phase boundaries for which $\xi - g$ is not zero in a neighbourhood of \mathfrak{Q}_i . Thus in the range of parameters r, t used in the expansion

(I.9) the field ϕ can be translated by ξ -g as before, with

$$\log \frac{\mathrm{d}\psi_{\mathrm{C}(\mathrm{t})}}{\mathrm{d}\phi_{\mathrm{C}(\mathrm{t})}} = \mathrm{F}(\Lambda, \Sigma)$$

given by (I.6) independent of t.

We now define the r, t-dependent irreducible kernels by the same Neumann series as in [K] in terms of C(t) and partially amputated Schwinger functions

$$S(\mathbf{x}_{1}, \dots, \mathbf{x}_{m}, y_{1}, \dots, y_{n})_{r,t} = \frac{\sum_{i \in \mathfrak{Y}} \prod_{i \in \mathfrak{Y}} r_{i}^{|\Sigma \cap N(\mathfrak{Q}_{i})|} \int d\phi_{C(t)} \prod_{i=1}^{m} \frac{d}{d\phi(x_{i})} \prod_{j=1}^{n} \phi(y_{j}) \chi_{\Sigma} e^{-V(\Lambda)}}{\sum_{\Sigma} \prod_{i \in \mathfrak{Y}} r_{i}^{|\Sigma \cap N(\mathfrak{Q}_{i})|} \int d\phi_{C(t)} \chi_{\Sigma} e^{-V(\Lambda)}}$$
(I.12)

They have the factorization property

$$S(\mathbf{x}_{1}, \dots, \mathbf{x}_{m+m'}, y_{1}, \dots, y_{n+n'})_{r,t} = S(\mathbf{x}_{1}, \dots, \mathbf{x}_{m}, y_{1}, \dots, y_{n})_{r,t} S(\mathbf{x}_{m+1}, \dots, \mathbf{x}_{m+m'}, y_{n+1}, \dots, y_{n+n'})_{r,t}$$
(I.13)

if \mathfrak{Q}_i separates $\{x_1, \ldots, x_m, y_1, \ldots, y_n\}$ from $\{x_{m+1}, \ldots, x_{m+m'}, y_{n+1}, \ldots, y_{n+n'}\}$ and $r_i = t_i = 0$. To calculate t-derivatives we use the following formula

$$\partial_{t_{i}} S(\mathbf{x}_{1}, \dots, \mathbf{x}_{m}, y_{1}, \dots, y_{n})_{r,t} = \frac{1}{2} \int d^{4}z S(\mathbf{z}_{1}, \mathbf{z}_{2}; \mathbf{x}_{1}, \dots, \mathbf{x}_{m}, y_{1}, \dots, y_{n})_{r,t} \partial_{t_{i}} C(t; z_{1}, z_{2})$$
(I.14)

where; denotes truncation. This formula is proved as usually by introducing cutoffs. It is now possible to repeat the same calculations as in [K, Section III]. This shows formally

Lemma 3. Let $x_1^0 < x_2^0 < y_1^0 < y_2^0$ and let $\alpha : \mathbb{Z} \to \{0, 1, 2\}$ be given such that $\alpha(i) \le 1$ if $x_1^0 < i < x_2^0$ or $y_1^0 < i < y_2^0$. Then

$$\begin{split} &\prod_{x_1^0 < i < x_2^0} \partial_{t_i}^{\alpha(i)} k(x_1; x_2)_{r,t} \left|_{t_i=0} = 0, \right. \\ &\prod_{x_1^0 < i < y_1^0} \partial_{t_i}^{\alpha(i)} L(x_1, x_2; y_1)_{r,t} \left|_{t_i=0} = 0 \end{split}$$

and

$$\prod_{\substack{k_1^0 < i < y_2^0}} \partial_{t_i}^{\alpha(i)} K(x_1, x_2; y_1, y_2)_{r,t} \big|_{t_i=0} = 0,$$

where $r_i\!=\!0$ whenever $t_i\!\neq\!1.$ Under the same condition on (r,t) and for $n\!=\!1,2,3$ we have

$$\prod_{\substack{\mathbf{x}_{k}^{0} < i < y_{1}^{0}}} \partial_{t_{i}}^{\alpha(i)} S_{k,l}^{n}(\mathbf{x}_{1}, \ldots, \mathbf{x}_{k}; \mathbf{y}_{1}, \ldots, \mathbf{y}_{l})_{\mathbf{r}, \mathbf{t}} \big|_{t_{i}=0} = 0$$

if $\alpha(i) < n$ and $\mathbf{x}_{1}^{0} < \cdots < \mathbf{x}_{k}^{0} < \mathbf{y}_{1}^{0} < \cdots < \mathbf{y}_{l}^{0}$.

Vol. 53, 1980 Particles exist in the low temperature φ_2^4 model

Suppose now that for a certain kernel $K_{r,t}(X; Y)$ we have shown that

$$\partial_t^{\alpha} K_{\mathbf{r},\mathbf{t}}(\mathbf{X};\mathbf{Y}) |_{\mathbf{r}^{\mathrm{I}}=\mathbf{t}^{\mathrm{I}}=\mathbf{0}} = 0$$

for every $I \subset \mathfrak{F}$ and every $\alpha: I \to \mathbb{Z}_+$ which is pointwise smaller than some $\beta = \beta(K)$. The r-t-expansion for $K(X; Y) = K_{1,1}(X; Y)$ is then

$$K(X, Y) = \sum_{\substack{I \cup J = \mathfrak{F} \\ I \cap J = \emptyset}} \int_{\{0\}}^{\{1\}} dr^{J} \int_{\{0\}}^{\{1\}} dt^{I} \prod_{i \in I} \frac{(1 - t_{i})^{\beta(i) - 1}}{(\beta(i) - 1)!} \, \partial_{r}^{J} \partial_{t}^{\beta} K_{r,t}(X; Y) \Big|_{\substack{r^{I} = 0 \\ t^{J} = 1}}$$
(I.15)

For the kernels we are interested in, \Im is given by the set of i's considered in the previous lemma, $\Im = l\mathbb{Z} \cap (x^0, y^0)$ for some interval $(x^0, y^0) \subset \mathbb{R}$. Then since .

$$\sum_{\substack{\mathbf{I} \cup \mathbf{J} = \mathfrak{F} \\ \mathbf{I} \cap \mathbf{J} = \emptyset}} 1 = 2^{|\mathfrak{F}|} \leq e^{l^{-1} |\mathbf{x}^0 - \mathbf{y}^0|}$$

the proof of Proposition 2 is reduced to the proof of the following lemma.

Lemma 4. Let $\varepsilon < 0$ be given and let (K(X; Y), O, \mathfrak{F}) denote one of the kernels considered in Proposition 2 and Lemma 3, with the corresponding order symbol O and index sets $\mathfrak{F} = \mathfrak{F}(K, X, Y)$. Then we have

$$\partial_t^{\alpha} \partial_r^J K_{r,t}(X;Y) \Big|_{r^{I}=0} = O(1) \exp\left(-(1-4\varepsilon) \sum_{i \in I} \alpha(i) - L |J|\right)$$
(I.16)

for $r = r^3$, $t = t^I$, $I, J \subset \mathfrak{F}$ and $\alpha : I \to \{0, 1, 2, 3\}$ depending on X, Y (taking constant values on products of unit lattice squares) such that $I \cap J = \emptyset$ and $r_i, t_i \in [0, 1]$.

The proof of this lemma will be given in the next section. It is based on a bound for modified expectations $\langle \Phi \rangle_{r,t,h}$ which will be shown in Section IV to be analytic in the variables r_i and $h(\alpha)$ in a large domain, and on a connection between t_i -derivatives and derivatives with respect to the new variables $h(\alpha)$. Corresponding irreducible kernels $K_{r,t,h}$ with $K_{r,t,0} = K_{r,t}$ are defined by convergent Neumann series in terms of expectations $\langle \Phi \rangle_{r,t,h}$ and thus Cauchy's formula can be applied to calculate $\partial_r^J \partial_t^\beta K_{r,t}$.

Remark. Recently J. Imbrie [I] has proven Theorem 1 by using a method which also works for small external field.

II. The expansion

In this section we define expectations $\langle \Phi \rangle_{r,t,h}$ and their expansion. We fix two sets I, J with $\mathfrak{F} = I \cup J \subset l\mathbb{Z} \cap \{x^0 : x \in \Lambda\}$, $I \cap J = \emptyset$. Then the α 's which occur in (I.16) are elements of I⁽⁴⁾, where I^(m) is the set of maps from I to $\{0, 1, \ldots, m-1\}$. Let

$$\partial_{t_i} C(t) \cdot \Delta_{\phi} = \frac{1}{2} \int d^2 x \, d^2 y \partial_{t_i} C(t; x, y) \frac{d^2}{d\phi(x) \, d\phi(y)}$$

Then multiple t-derivatives can be written as

$$\partial_{t}^{\beta} \int d\phi_{C(t)} Q = \sum_{\substack{\alpha_{1} + \dots + \alpha_{l} = \beta \\ \alpha_{j} \leq 1}} \int d\phi_{C(t)} \prod_{j=1}^{l} \left[\partial_{t}^{\alpha_{j}} C(t) \cdot \Delta_{\phi} \right] Q \tag{II.1}$$

or, by introducing variables $h(\alpha_i)$

$$\partial_{t}^{\beta} \int d\phi_{C(t)} Q$$

$$= \sum_{\substack{\alpha_{1}+\dots+\alpha_{l}=\beta \\ \alpha_{i}\leq 1}} \prod_{i=1}^{l} \partial_{h(\alpha_{i})} \int d\phi_{C(t)} \prod_{\alpha\in\pi_{0}} \prod_{j\in \mathbb{Z}^{4}} [1+h(\alpha)\partial_{t}^{\alpha}C_{j}(t)\cdot\Delta_{\phi}] Q|_{h=0}$$
(II.2)

where $C_j(t; x, y) = \Delta_{j1}(x)C(t; x, y)\Delta_{j2}(y)$ for $j = (j_1, j_2) \in l\mathbb{Z}^4$, and Δ_i denotes the square $[i^0, i^0 + 1] \times [i^1, i^1 + 1]$, respectively its characteristic function. π_0 is a subset of $P(I^{(4)})$, $P(I^{(m)}) = \{\pi \subset I^{(m)} : \sum_{\alpha \in \pi} \alpha \in I^{(m)}, \emptyset \notin \pi\}$, which contains $\alpha_1, \ldots, \alpha_l$.

We now define

$$\langle \Phi \rangle_{\mathbf{r},\mathbf{t},\mathbf{h}}$$

$$= \frac{\sum_{\Sigma} \prod_{i \in \mathfrak{Y}} r_{i}^{|\Sigma \cap \mathbf{N}(\mathfrak{L}_{i})|} \int d\phi_{\mathbf{C}(\mathbf{t})} \prod_{\alpha \in \pi_{0}} \prod_{j \in l\mathbb{Z}^{4}} [1 + h(\alpha)\partial_{\mathbf{t}}^{\alpha}\mathbf{C}_{j}(\mathbf{t}) \cdot \Delta_{\phi}] \Phi \chi_{\Sigma} e^{-\mathbf{V}(\Lambda)}}{\sum_{\Sigma} \prod_{i \in \mathfrak{Y}} r_{i}^{|\Sigma \cap \mathbf{N}(\mathfrak{L}_{i})|} \int d\phi_{\mathbf{C}(\mathbf{t})} \prod_{\alpha \in \pi_{0}} \prod_{j \in l\mathbb{Z}^{4}} [1 + h(\alpha)\partial_{\mathbf{t}}^{\alpha}\mathbf{C}_{j}(\mathbf{t}) \cdot \Delta_{\phi}] \chi_{\Sigma} e^{-\mathbf{V}(\Lambda)}}$$
(II.3)

The set $\pi_0 \in P(I^{(4)})$ is always supposed to be chosen according to the derivatives that we want to compute, and the domain of parameters (r, h) which we consider is (with ε given by Lemma 4)

 $\mathfrak{R} = \{ (\mathbf{r}, \mathbf{h}) : \mathbf{r} \in \mathbb{C}^{|\mathfrak{V}|}, \, \mathbf{h} \in \mathbb{C}^{|\pi_0|}, \, \mathbf{r}_i = 0 \text{ for } i \in \mathbf{I},$

$$|\mathbf{r}_i| < e^{3L}$$
 for $i \in \mathfrak{F}$, $|\mathbf{h}(\alpha)| < \exp((1-2\varepsilon)(\mathbf{d}(\alpha)+l))$ for $\alpha \in \pi_0$ }

where $d(\alpha) = \max\{|i-j|: \alpha(i), \alpha(j) \neq 0\}$ for $\alpha \in I^{(2)}$.

To construct our kernels K_{r,t,h} we need fields of the form

$$\Phi = \prod_{i=1}^{M} \Box_{j_i}(x_i) \frac{d}{d\phi(x_i)} \prod_{i=M+1}^{N} \Box_{j_i}(x_i) : \phi^{\nu_i} : (x_i)$$
(II.4)

where \Box_j denotes the unit square $[j^0, j^0+1] \times [j^1, j^1+1]$, respectively its characteristic function. The support $\Omega = \Box_{j1} \cup \cdots \cup \Box_{j_N}$ is supposed to be in Λ . We define two functions m_{Φ} and n_{Φ} by

$$\begin{split} \mathbf{m}_{\Phi}(\Omega') &= \sum_{i=1}^{M} |\Omega' \cap \Box_{\mathbf{j}_{i}}|, \\ \mathbf{n}_{\Phi}(\Omega') &= \sum_{i=M+1}^{N} |\Omega' \cap \Box_{\mathbf{j}_{i}}| \cdot \mathbf{\nu}_{i}, \end{split}$$

for $\Omega' \subset \mathbb{R}^2$. The expectation is considered as a function of the variables x_1, \ldots, x_N . Since $d/d\phi(x_i)$ can contract to some $: \phi^{\nu_i}: (x_i)$ or to a vertex V together

with some other $d/d\phi(x_i)$, this function has the form

$$\mathbf{f}(\mathbf{x}_1,\ldots,\mathbf{x}_N) = \sum_{\pi \in \mathbf{P}(\{1,\ldots,N\})} \int \prod_{j=1}^{|\pi|} \mathrm{d}^2 \mathbf{y}_j \prod_{i \in \sigma_j} \delta(\mathbf{x}_i - \mathbf{y}_i) \cdot \mathbf{f}_{\pi}(\mathbf{y}_1,\ldots,\mathbf{y}_{|\pi|})$$

where the sum runs over partitions $\pi = \{\sigma_1, \ldots, \sigma_{|\pi|}\}$ of $\{1, \ldots, N\}$. Let π denote a fixed partition of $\{1, \ldots, N_0\}$, $N_0 \ge N$. We define

$$|\mathbf{f}|_{\mathbf{q}} = \|\mathbf{f}_{\pi_1}\|_{\mathbf{L}\mathbf{q}}$$

if $\pi = \pi_1 \cup \pi_2$ for some partition π_1 of $\{1, \ldots, N\}$. Otherwise let $|f|_q = 0$. Notice that $|f|_q = |f^{(1)}|_q \cdot |f^{(2)}|_q$ if $f = f_1 \otimes f_2$.

Lemma 5. Fix $q \ge 1$ and $\varepsilon < 0$. Then there are positive constants β and ρ such that for $(r, h) \in \Re$, $t \in [0, 1]^I$ and for $\lambda > 0$ sufficiently small $\langle \Phi \rangle_{r,t,h}$ is analytic in r and h. Furthermore for $m_{\Phi} + n_{\Phi} \neq 0$

$$|\langle \Phi \rangle_{\mathbf{r},\mathbf{t},\mathbf{h}}|_{\mathbf{q}} \leq C_1(\mathbf{m}_{\Phi},\mathbf{q})C_2(\mathbf{n}_{\Phi},\mathbf{q})$$

independently of Λ , where $C_1(0, q) = 1$, $C_2(0, q) = \lambda^{1/2-\epsilon}$ and for m, $n \neq 0$

$$C_{1}(m, q) = e^{\rho m(\Omega)^{4/3}}$$

$$C_{2}(n, q) = (1 + \lambda^{-n(\Omega)/2} e^{-\beta \lambda^{-1/2}}) e^{\ln(\Omega)} \prod_{\Box \subset \Lambda} n(\Box)!.$$

This lemma will be proven in the next section, together with the following one. Let

$$\Phi_{\{i\}} = \Box_{j_i}(\mathbf{x}_i) \frac{\mathrm{d}}{\mathrm{d}\phi(\mathbf{x}_i)} \qquad i = 1, \dots, M,$$

$$\Phi_{\{i\}} = \Box_{j_i}(\mathbf{x}_i) : \phi^{\nu_i} : (\mathbf{x}_i) \qquad i = M+1, \dots, N$$

be given with $\Omega = \Box_{j_1} \cup \cdots \cup \Box_{j_N} \subset \Lambda$. To subsets $\Omega' \subset \Omega$ and $R' \subset R = \{1, \ldots, N\}$ we associate products

$$\Phi_{\Omega'} = \prod_{\substack{i \le M \\ \Box_{i_i} \subset \Omega'}} \Phi_{\{i\}} \prod_{\substack{i > M \\ \Box_{i_i} \subset \Omega'}} \Phi_{\{i\}}$$

$$\Phi(\mathbf{R}') = \prod_{\substack{i \le M \\ i \in \mathbf{R}'}} \Phi_{\{i\}} \prod_{\substack{i > M \\ i \in \mathbf{R}'}} \Phi_{\{i\}}.$$
(II.5)

Then for partitions $\{R_1,\ldots,\,R_s\}$ of R we define truncated expectations as usually by

$$\langle \Phi(\mathbf{R}_1); \Phi(\mathbf{R}_2); \dots; \Phi(\mathbf{R}_s) \rangle = -\sum_{p=1}^{s} \frac{(-1)^p}{p} \sum_{\sigma_1 \cup \dots \cup \sigma_p = \{1, \dots, s\}} \prod_{j=1}^{p} \left\langle \Phi\left(\bigcup_{i \in \sigma_j} \mathbf{R}_i\right) \right\rangle$$
(II.6)

Lemma 6. For given $q \ge 1$ there are constants K, $\delta > 0$ such that for $(r, h) \in \Re$, $t \in [0, 1]^{I}$ and for $\lambda > 0$ sufficiently small

 $\left| \langle \Phi(\mathbf{R}_1); \ldots; \Phi(\mathbf{R}_s) \rangle_{\mathbf{r},\mathbf{t},\mathbf{h}} \right|_{\mathbf{q}} \leq (\mathbf{K}s)^s C_1(\mathbf{m}_{\Phi(\mathbf{R})}, \mathbf{q}) C_2(\mathbf{n}_{\Phi(\mathbf{R})}, \mathbf{q}) e^{-\delta d(\Omega_1, \ldots, \Omega_s)}$

where Ω_k is the support of $\Phi(R_k)$ and $d(\Omega_1, \ldots, \Omega_s) = \sup \{|x-y| : x \in \Omega_i, y \in \Omega_j, 1 \le i < j \le s\}$.

Proof of Lemma 3 and 4. With (I.14) and Lemma 5 and 6 as input we can exactly follows [K] in order to show that irreducible kernels $K_{r,t,h}$ are well defined (i.e. that Proposition 2 holds for $1-\varepsilon > 0$ sufficiently small), that they are analytic in r and h for $(r, h) \in \Re$, and that they verify the statements of Lemma 3.

Let now K denote one of the kernels k, L, K or $S_{k,l}^n$. Then from (II.2) it follows [S III] that

$$\partial_{t}^{\alpha} K_{r,t} = \sum_{\substack{\alpha_{1} + \dots + \alpha_{j} = \alpha \\ \alpha_{i} \leq 1}} \prod_{i=1}^{l} \partial_{h(\alpha_{i})} K_{r,t,h} \Big|_{h=0}$$
(II.7)

and thus

$$\partial_{t}^{\alpha}\partial_{r}^{J}K(X;Y)_{r,t} = \sum_{\substack{\alpha_{1}+\dots+\alpha_{l}=\alpha\\\alpha_{j}\leq 1}} \oint \prod_{\alpha\in\{\alpha_{1},\dots,\alpha_{l}\}} \frac{h(\alpha)! \, dh(\alpha)}{2\pi i h(\alpha)^{n(\alpha)+1}} \oint \prod_{j\in J} \frac{dz_{j}}{2\pi i (z_{j}-r_{j})^{2}} K(X;Y)_{z,t,h}$$
(II.8)

where $n(\alpha)$ is the number of copies of α in $(\alpha_1, \ldots, \alpha_l)$ and where the integrations are over circles

$$|\mathbf{h}(\alpha)| = \exp\left((1 - 3\varepsilon)(\mathbf{d}(\alpha) + l)\right),$$

 $|\mathbf{z}_i| = e^{2L}.$

We suppose r_i , $t_i \in [0, 1]$, $r^I = 0$. Since $n(\alpha) \le 3$ and since the number of sets $(\alpha_1, \ldots, \alpha_k)$ with $0 \ne \alpha_j \le 1$, $\sum_{j=1}^k \alpha_j = \alpha$ and $\sum_{j=1}^k (d(\alpha_j) + l) = ml$ is zero for $m < |\alpha| = \sum_{i \in I} \alpha(i)$ and bounded by 32^m for $m \ge |\alpha|$, we obtain

$$\partial_{t}^{\alpha}\partial_{r}^{J}K(X;Y)_{r,t} = O(1) \cdot \sum_{m \ge |\alpha|} 32^{m} \cdot 6^{m} \exp\left(-(1-3\varepsilon)ml - L |J|\right)$$
$$= O(1) \exp\left(-(1-4\varepsilon)l |\alpha| - L |J|\right).$$

Following [GJS III] we now define a cluster expansion for expectations $F_{\Omega \cap X, X, \Sigma}(r, t, h, s = 1)$, which will be defined below so that for $r^{I} = 0$

$$\langle \Phi \rangle_{\mathbf{r},\mathbf{t},\mathbf{h}} = \frac{S_{\Omega}}{S_{\emptyset}} = \frac{\sum_{\Sigma} F_{\Omega,\mathbf{Y},\Sigma}(\mathbf{r},\mathbf{t},\mathbf{h},\mathbf{s}=1)}{\sum_{\Sigma} F_{\emptyset,\mathbf{Y},\Sigma}(\mathbf{r},\mathbf{t},\mathbf{h},\mathbf{s}=1)}$$
(II.9)

with $\mathbf{Y} = (\mathbf{Y}, \partial \mathbf{Y}, \emptyset)$. Φ denotes a fixed field of the form (II.4).

Let X be a subset of Y built up from lattice squares Δ and let $\partial X = \partial X^+ \cup \partial X^$ be a partition of its boundary. We suppose the triple $\mathbf{X} = (X, \partial X^+, \partial X^-)$ to be regular, i.e.

 $N(\partial X^{-}) \cap (N(\partial X^{+}) \cup (Y \setminus \Lambda)) = \emptyset.$

For each $X' \subset X$ let b(X') denote the set of $(lZ)^2$ lattice bonds b in X' with $|b \cap \partial X'| = 0$. Furthermore let $A(\mathbf{X})$ denote the set of spin configurations Σ on X such that $\Sigma(\Box) = \pm$ for $\Box \subset N(\partial X^{\pm})$. Given a configuration $\Sigma \in A(\mathbf{X})$ we introduce a cluster expansion which only involves bounds b far from phase boundaries, i.e. $b \in b = b(X \setminus N(\Sigma))$.

Let C(t, s) be the covariance

$$C(t, s) = \prod_{b \in b} [(1 - s_b)\theta_b + s_b] \prod_{i \in I} [(1 + t_i)\theta_{\mathfrak{L}_i} + t_i] (-\Delta_{\partial x} + 1)^{-1}$$
(II.10)

for $t \in [0, 1]^{I}$, $s \in [0, 1]^{b}$, where $\theta_{\Gamma_{2}}(-\Delta_{\Gamma_{1}}+1)^{-1} = (-\Delta_{\Gamma_{1}\cup\Gamma_{2}}+1)^{-1}$, and where Δ_{Γ} denotes the Laplacean with zero Dirichlet data on each $b \in \Gamma$. For $\Gamma \subset b$ and for F = F(s) we define $\Gamma^{c} = b \setminus \Gamma$, $\partial_{s}^{\Gamma} = \prod_{b \in \Gamma} \partial_{s_{b}}$,

$$\mathbf{s}(\Gamma^{\mathbf{c}})_{\mathbf{b}} = \begin{cases} 0 & \text{if } \mathbf{b} \in \Gamma^{\mathbf{c}} \\ \mathbf{s}_{\mathbf{b}} & \text{if } \mathbf{b} \in \Gamma \end{cases}$$

and

$$\delta_{s}^{\Gamma}F = \int_{\{0\}}^{\{1\}} ds \partial_{s}^{\Gamma}F(s(\Gamma^{c})).$$

The cluster expansion is then the identity

$$\mathbf{F}(1) = \sum_{\Gamma \subset b} \delta_{s}^{\Gamma} \mathbf{F}, \qquad (II.11)$$

applied to expectations $(\Omega' \subset X)$

$$F_{\Omega',\mathbf{X},\Sigma}(\mathbf{r},\mathbf{t},\mathbf{h},\mathbf{s}) = \prod_{i\in\mathfrak{Y}} r_i^{|\Sigma\cap N(\mathfrak{Q}_i)|} \int d\Psi_{C(\mathbf{t},\mathbf{s})} e^{-F(\Lambda,\Sigma)}$$

$$\times \prod_{\alpha\in\pi_0} \prod_{j\in l\mathbb{Z}^4} [1+h(\alpha)\partial_t^{\alpha}C_j(\mathbf{t},\mathbf{s})\cdot\Delta_{\phi}] \Phi_{\Omega'}\chi_{\Sigma} e^{-\mathbf{V}(\Lambda)}.$$
(II.12)

We continue by considering $\mathbf{X} = \mathbf{Y}$. Equation (II.9) is obtained from (II.3) through the translation $\phi \mapsto \psi = \phi + \xi - g$ (see (I.6)) by using that

$$\frac{\mathrm{d}\psi_{\mathrm{C}(\mathrm{t},\mathrm{s})}}{\mathrm{d}\phi_{\mathrm{C}(\mathrm{t},\mathrm{s})}} = \frac{\mathrm{d}\psi_{\mathrm{C}(1,1)}}{\mathrm{d}\phi_{\mathrm{C}(1,1)}} = \mathrm{e}^{-\mathrm{F}(\Lambda,\Sigma)}$$

for our particular choice of Dirichlet lines.

In the usual way we order the sum (II.11) such that subsums are indexed by a covering of Y with closed, nonempty connected sets X_1, \ldots, X_n having pairwise disjoint interior. The set of such coverings will be denoted by C(Y). For given $\{X_1, \ldots, X_n\} \in C(Y)$ we sum first over those Γ 's satisfying

$$Y \setminus \left(\partial Y \bigcup_{b \in \Gamma^{c}} b\right) = \bigcup_{j=1}^{n} Y_{j}$$
(II.13)

with Y_j open, connected and closure $(Y_j) = X_j$ for j = 1, ..., n. Of course the X_j must be compatible with Σ , i.e. $\partial X_j \subset b(Y \setminus N(\Sigma))$. If we orient their boundaries by defining

$$\partial \mathbf{X}_{i}^{\pm} = \partial \mathbf{X}_{i} \cap \{\Box : \Sigma(\Box) = \pm\},\$$

then they are pairwise compatible, i.e.

$$N(\partial X_i^-) \cap ((N(\partial X_i^+) \cup (Y \setminus \Lambda)) = \emptyset.$$
(II.14)

The restrictions of Σ and Γ to X_j will be denoted by Σ_j and Γ_j respectively. Notice that $\Sigma_j \in A(\mathbf{X}_j)$ and $\Gamma_j \in B(X_j, \Sigma_j)$, where

$$\mathbf{B}(\mathbf{X}, \boldsymbol{\Sigma}) = \bigg\{ \Gamma \in b(\mathbf{X} \setminus \mathbf{N}(\boldsymbol{\Sigma})) : \mathbf{X} \setminus \bigcup_{\mathbf{b} \in \Gamma^c} \mathbf{b} \text{ is connected} \bigg\}.$$

Furthermore we have the factorization property

$$F_{\Omega,\mathbf{Y},\boldsymbol{\Sigma}}(\mathbf{r},\mathbf{t},\mathbf{h},\mathbf{s}(\Gamma^{c})) = \prod_{j=1}^{n} F_{\Omega\cap \mathbf{X}_{j},\mathbf{X}_{j},\boldsymbol{\Sigma}_{j}}(\mathbf{r},\mathbf{t},\mathbf{h},\mathbf{s}(\Gamma_{j}^{c})).$$
(II.15)

.

The expansion for S_{Ω} and S_{\emptyset} can now be written as follows. Let $\Omega' \subset \Omega$ be given and let $C(\mathbf{Y})$ denote the set of sets $\{\mathbf{X}_1, \ldots, \mathbf{X}_n\}$ with pairwise compatible elements, such that $\{X_1, \ldots, X_n\} \in C(Y)$. Then

$$\begin{split} \mathbf{S}_{\Omega'} &= \sum_{\boldsymbol{\Sigma}} \sum_{\Gamma \in \mathbf{b}(\mathbf{Y})} \delta_{s}^{\Gamma} \mathbf{F}_{\Omega', \mathbf{Y}, \boldsymbol{\Sigma}} \\ &= \sum_{\boldsymbol{\Sigma}} \sum_{\Gamma \in \mathbf{b}(\mathbf{Y} \setminus \mathbf{N}(\boldsymbol{\Sigma}))} \prod_{j} \delta_{s}^{\Gamma_{j}} \mathbf{F}_{\Omega' \cap \mathbf{X}_{j}, \mathbf{X}_{j}, \boldsymbol{\Sigma}_{j}} \\ &= \sum_{\{\mathbf{X}_{j}\} \in \mathbf{C}(\mathbf{Y})} \prod_{j} \sum_{\boldsymbol{\Sigma} \in \mathbf{A}(\mathbf{X}_{j})} \sum_{\Gamma \in \mathbf{B}(\mathbf{X}_{j}, \boldsymbol{\Sigma})} \delta_{s}^{\Gamma} \mathbf{F}_{\Omega' \cap \mathbf{X}_{j}, \mathbf{X}_{j}, \boldsymbol{\Sigma}}. \end{split}$$
(II.16)

The proof of the convergence of the expansion for $\langle \Phi \rangle = S_{\Omega} S_{\emptyset}^{-1}$ can now be based on the following estimate.

Lemma 7. Fix $q \ge 1$ and $\varepsilon > 0$. Then there are constants a, b, c > 0 such that for regular $\mathbf{X}, \Omega' \subset \Omega \cap \mathbf{X}$, $(\mathbf{r}, \mathbf{h}) \in \Re$, $\mathbf{t} \in [0, 1]^{\mathrm{I}}$ and for $\lambda > 0$ sufficiently small

$$\begin{aligned} |\delta_s^1 F_{\Omega', \mathbf{X}, \Sigma}(\mathbf{r}, \mathbf{t}, \mathbf{h}, \mathbf{s})|_q &\leq C_1(m_{\Phi_{\Omega'}}, q) C_3(\mathbf{n}_{\Phi}, \Omega', q) \\ &\times \exp\left(-c \left|\Gamma\right| - b\lambda^{-1/2} \left|\mathbf{N}(\Sigma) \cap \mathbf{X}\right| + al^{-2} \left|\mathbf{X}\right|\right) \end{aligned}$$

where

$$C_3(n, \Omega', q) = C(\mathbf{X}, \Sigma)^{n(\Omega')} \exp\left(\frac{1}{2} + 8c\right) \ln\left(\Omega'\right) \prod_{\Box \subset \Omega'} n(\Box)!$$

if $n \neq 0$, $\Omega' \neq \emptyset$,

$$C(\mathbf{X}, \Sigma) = \begin{cases} 1 & if \quad |\Sigma| = |\partial X^{-}| = 0\\ 2\xi & otherwise \end{cases},$$
$$C_{3}(0, \Omega', q) = \begin{cases} 1 & if \quad \Omega' = \emptyset\\ \lambda^{1/2-\varepsilon} & otherwise \end{cases}.$$

This lemma will be proved in the last section.

III. Convergence of the expansion

In this section we prove the convergence of the low temperature cluster expansion for the modified expectations

$$\langle \Phi \rangle_{\mathbf{r},\mathbf{t},\mathbf{h}} = \frac{\sum_{\{\mathbf{Y}_{j}\} \in \mathbf{C}(\mathbf{Y})} \prod_{j} \sum_{\Sigma \in \mathbf{A}(\mathbf{Y}_{j})} \sum_{\Gamma \in \mathbf{B}(\mathbf{Y}_{j},\Sigma)} \delta_{s}^{\Gamma} F_{\Omega' \cap \mathbf{Y}_{j},\mathbf{Y},\Sigma}(\mathbf{r},\mathbf{t},\mathbf{h},s)}{\sum_{\{\mathbf{Y}_{j}\} \in \mathbf{C}(\mathbf{Y})} \prod_{j} \sum_{\Sigma \in \mathbf{A}(\mathbf{Y}_{j})} \sum_{\Gamma \in \mathbf{B}(\mathbf{Y}_{j},\Sigma)} \delta_{s}^{\Gamma} F_{\emptyset,\mathbf{Y}_{j},\Sigma}(\mathbf{r},\mathbf{t},\mathbf{h},s)}$$
(III.1)

and the exponential clustering of the corresponding truncated expectations. We use a method due to H. Kunz, B. Souillard, T. Balaban and K. Gawedski [BG].

First the numerator and denominator in (III.1) are divided by a product $\prod_{\Delta \subset \Lambda} Z_{\Delta}$ of $l\mathbb{Z}^2$ square partition functions

$$Z_{\Delta} = \sum_{\Gamma \in \mathbf{B}(\Delta, \emptyset)} \delta_{s}^{\Gamma} F_{\theta, \Delta, \theta}(0, 0, 0, s)$$

which do not depend on the sign of the boundary $\partial \Delta$. The expansion can then be written in terms of the following quantities.

$$\mathbf{S}_{\boldsymbol{\Omega}',\mathbf{X}}(\mathbf{r},\mathbf{t},\mathbf{h}) = \mathbf{Z}_{\Delta}^{-l^{-2}|\mathbf{X}\cap\Lambda|} \sum_{\boldsymbol{\Sigma}\in\mathbf{A}(\mathbf{X})} \sum_{\boldsymbol{\Gamma}\in\mathbf{B}(\mathbf{X},\boldsymbol{\Sigma})} \delta_{s}^{\boldsymbol{\Gamma}} \mathbf{F}_{\boldsymbol{\Omega}',\mathbf{X},\boldsymbol{\Sigma}}(\mathbf{r},\mathbf{t},\mathbf{h},s)$$
(III.2)

if $\Omega' \subset \Omega \cap X$, $|X| > l^2$, X is connected and **X** regular; otherwise let $S_{\Omega', \mathbf{X}} = 0$. The new expression for $\langle \Phi \rangle_{r,t,h}$ is

$$\langle \Phi \rangle_{\mathbf{r},\mathbf{t},\mathbf{h}} = \frac{\sum_{j} \prod_{j} S_{\Omega \cap \mathbf{Y}_{j},\mathbf{Y}_{j}}}{\sum_{\{\mathbf{Y}_{j}\} \text{ comp. } j} \prod_{j} S_{\phi,\mathbf{Y}_{j}}}$$
(III.3)

The sum ranges over sets $M = \{Y_1, \ldots, Y_m\}$ such that $M \cup N$ is a partition of **Y** for some $N = \{X_1, \ldots, X_n\}$ with $|X_j| = l^2$. Such sets M will be called compatible.

Lemma 8. Fix $q \ge 1$. Then there is a positive constant δ such that for $X \subseteq Y$, $\Omega' \subseteq \Omega \cap X$, $(r, h) \in \Re$, $t_i \in [0, 1]$ and $\lambda > 0$ sufficiently small

$$\begin{split} &|S_{\emptyset,\mathbf{X}}(\mathbf{r},t,h)| \leq \! e^{-2\delta|X|l^{-1}} \\ &|S_{\Omega',\mathbf{X}}(\mathbf{r},t,h)|_{q} \! \leq \! C_{1}(m_{\Phi_{\Omega'}},q) C_{2}(n_{\Phi_{\Omega'}},q) e^{-2\delta|X|l^{-1}}. \end{split}$$

Proof. The proof is as in [GJS III]: One part of the factor $e^{-b\lambda^{-1/2}|N(\Sigma)|}$ in Lemma 7 is used for the first factor in C₂ (see Lemma 5). The other part controls the first sum in (III.2).

$$\sum_{\boldsymbol{\Sigma} \in A(\mathbf{X})} \exp\left(-b'\lambda^{-1/2} |\mathbf{N}(\boldsymbol{\Sigma})|\right) \leq \sum_{\boldsymbol{\Sigma} \in A(\mathbf{X})} \exp\left(-\frac{1}{2}b'\lambda^{-1/2} |\boldsymbol{\Sigma}|\right)$$

$$\leq (1 + \exp\left(-\frac{1}{2}b'\lambda^{-1/2}\right))^{2|\mathbf{X}|} \leq e^{\lambda|\mathbf{X}|}.$$
(III.5)

To bound the second sum in (III.2) notice that for $\Gamma \in B(X, \Sigma)$ we have $2l^{-1} |\Gamma| \ge |X| l^{-2} - 1$. Thus

$$\sum_{\Gamma \in B(\mathbf{X}, \Sigma)} e^{-c|\Gamma|} \le \exp\left(-\frac{c}{2} l^{-1}(|\mathbf{X}| - l^2)\right) \sum_{\Gamma \in B(\mathbf{X}, \Sigma)} 1$$

$$\le \exp\left(-\frac{c}{2} l^{-1}(|\mathbf{X}| - l^2)\right) 2^{2l^{-2}|\mathbf{X}| - (1/2)l^{-1}|\partial \mathbf{X}|} \le \exp\left(-\frac{c}{4} l^{-1}(|\mathbf{X}| - l^2)\right),$$
(III.6)

and we can take $2\delta = c/10$ if we put a factor $e^{\lambda l^2 + 2\delta l}$ into C_2 for the case $|X| = l^2$. Finally the factor $Z_{\Delta}^{-l^{-2}|X \cap \Lambda|}$ is bounded by using [GJS III]

$$|Z_{\Delta}^{\pm 1}| \le e^{O(1)\lambda^{1/2}l^2}.$$
 (III.7)

We continue by expanding the ratio on the left hand side of (III.3). Notice that a compatible set $\{\mathbf{Y}_1, \ldots, \mathbf{Y}_m\}$ is already determined by $\{\mathbf{Y}_1, \ldots, \mathbf{Y}_m\}$ and the relative signs of the boundary loops (since the outermost loop must have the + sign), i.e. by $\{\tilde{\mathbf{Y}}_1, \ldots, \tilde{\mathbf{Y}}_m\}$, where $\tilde{\mathbf{X}} = \{(\mathbf{X}, \partial \mathbf{X}^+, \partial \mathbf{X}^-), (\mathbf{X}, \partial \mathbf{X}^-, \partial \mathbf{X}^+)\}$. The starting point for the expansion of [BG] is the observation that the factor

$$\varphi = \begin{cases} \prod_{j=1}^{m} \mathbf{S}_{\mathbf{\Omega}_{j}, \mathbf{Y}_{j}} & \text{if } \{\mathbf{Y}_{1}, \dots, \mathbf{X}_{m}\} \text{ are compatible} \\ 0 & \text{otherwise} \end{cases}$$

can be expressed in terms of an 'interaction' $U(\alpha)$ between pairs $\alpha = ((\Omega_i, \tilde{\mathbf{Y}}_i), (\Omega_i, \tilde{\mathbf{Y}}_i))$.

Let us call \mathbf{X} positive or negative according to the sign of its external boundary loop. A negative (positive) boundary loop of a positive (negative) \mathbf{X} will be called inner. Then by regarding

$$S_m((\Omega_1, \mathbf{Y}_1), \ldots, (\Omega_m, \mathbf{Y}_m)) = S_{\Omega_1, \mathbf{Y}_1} \cdots S_{\Omega_m, \mathbf{Y}_m}$$

for fixed Ω_i , $\tilde{\mathbf{Y}}_i$, i = 1, ..., m as a function of m signs we define

$$\begin{bmatrix} U((\Omega_{1}, \tilde{\mathbf{Y}}_{1}), (\Omega_{2}, \tilde{\mathbf{Y}}_{2}))S_{m}]((\Omega_{1}, \mathbf{Y}_{1}), (\Omega_{2}, \mathbf{Y}_{2}), \ldots) \\ 0 \quad \text{if} \quad Y_{1} \cap Y_{2} \neq \emptyset. \\ S_{m}((\Omega_{1}, (Y_{1}, \partial Y_{1}^{-}, \partial Y_{1}^{+})), (\Omega_{2}, \mathbf{Y}_{2}), \ldots) \quad \text{if} \quad \Omega_{1} \neq \emptyset \\ \text{and if } Y_{1} \text{ is surrounded by an inner boundary} \\ \text{loop of } \mathbf{Y}_{2}. \\ S_{m}((\Omega_{1}, \mathbf{Y}_{1}), (\Omega_{2}, (Y_{2}, \partial Y_{2}^{-}, \partial Y_{2}^{+})), \ldots) \quad \text{if} \quad \Omega_{2} \neq \emptyset \\ \text{and if } Y_{2} \text{ is surrounded by an inner boundary} \\ \text{loop of } \mathbf{Y}_{1}. \\ S_{m}((\Omega_{1}, \mathbf{Y}_{1}), (\Omega_{2}, \mathbf{Y}_{2}), \ldots) \quad \text{otherwise.} \end{aligned}$$

The remaining $U(\alpha)$ for $\alpha \in \mathfrak{A} = \{((\Omega_i, \tilde{\mathbf{Y}}_i), (\Omega_j, \tilde{\mathbf{Y}}_j)): 1 \le i < j \le m\}$ are defined analogously. By identifying $\tilde{\mathbf{X}}$ with its positive element, φ can be written as follows

$$\varphi = \varphi((\Omega_1, \tilde{\mathbf{Y}}_1), \dots, (\Omega_m, \tilde{\mathbf{Y}}_m))$$

= $\prod_{\alpha \in \mathfrak{A}} U(\alpha) S_m((\Omega_1, \tilde{\mathbf{Y}}_1), \dots, (\Omega_m, \tilde{\mathbf{Y}}_m)).$ (III.9)

Here we have used that by the $\varphi \mapsto -\varphi$ symmetry $S_{\emptyset,(X,\partial X^+,\partial X^-)} = S_{\emptyset,(X,\partial X^-,\partial X^+)}$. Next we write a graphical expansion for $\prod_{\alpha \in \mathfrak{A}} U(\alpha)$, namely

$$\prod_{\alpha \in \mathfrak{A}} \mathbf{U}(\alpha) = \prod_{\Omega_i, \Omega_j \neq 0} \mathbf{U}((\Omega_i, \tilde{\mathbf{Y}}_i), (\Omega_j, \tilde{\mathbf{Y}}_j)) \sum_{\mathbf{G}_c} \prod_{\alpha \in \mathbf{G}_c} \mathbf{A}(\alpha) \prod_{\alpha \in \mathfrak{A}'} \mathbf{U}(\alpha)$$
(III.10)

where $A(\alpha) = U(\alpha) - 1$. The sum ranges over connected graphs G_c whose lines are a subset of those $\alpha = ((\Omega_i, \tilde{\mathbf{Y}}_i), (\Omega_j, \tilde{\mathbf{Y}}_j)) \in \mathfrak{A}$ with $\Omega_i = \emptyset$ or $\Omega_j = \emptyset$, and whose connected components contain at least one $(\Omega_j, \tilde{\mathbf{Y}}_j)$ with $\Omega_j \neq \emptyset$. \mathfrak{A}' contains the pairs of those points $\{(\emptyset, \mathbf{Y}_i)\}_{i \in I}$ which are not involved in G_c .

With a sum over the $\Omega_j \in \{\Omega \cap Y_j, \emptyset\}$ such that $\bigcup_j \Omega_j = \Omega$, the numerator in (III.3) can then be written as

$$\begin{split} &\sum_{m} \frac{1}{m!} \sum_{\substack{\mathbf{Y}_{1}, \dots, \mathbf{Y}_{m} \text{ pos.} \\ \Omega_{1}, \dots, \Omega_{m}}} \varphi((\Omega_{1}, \mathbf{Y}_{1}), \dots, (\Omega_{m}, \mathbf{Y}_{m})) \\ &= \sum_{m} \frac{1}{m!} \sum_{\substack{\mathbf{Y}_{1}, \dots, \mathbf{Y}_{m} \text{ pos. } \Omega_{i}, \Omega_{j} \neq \emptyset}} \prod_{\substack{\mathbf{U}((\Omega_{i}, \tilde{\mathbf{Y}}_{i}), (\Omega_{j}, \tilde{\mathbf{Y}}_{j})) \\ \Omega_{1}, \dots, \Omega_{m}}} \prod_{\boldsymbol{X}_{1}, \dots, \Omega_{m}} U((\Omega_{i}, \tilde{\mathbf{Y}}_{i}), (\Omega_{j}, \tilde{\mathbf{Y}}_{j})) \\ &\times \sum_{\mathbf{G}_{n}} \prod_{\alpha \in \mathbf{G}_{n}} \mathbf{A}(\alpha) \mathbf{S}_{m-|\mathbf{I}|}(\{(\Omega_{j}, \mathbf{Y}_{j})\}_{j \notin \mathbf{I}}) \prod_{\alpha \in \mathfrak{A}'} U(\alpha) \mathbf{S}_{|\mathbf{I}|}(\{(\emptyset, \mathbf{Y}_{i})\}_{i \in \mathbf{I}}) \end{split}$$
(III.11)

or, by summing first the terms with fixed $\{(\Omega_j, \mathbf{Y}_j)\}_{j \notin I} = \{(\Omega'_1, \mathbf{X}_1), \dots, (\Omega'_k, \mathbf{X}_k)\}, k = m - |I|,$

$$\sum_{m} \frac{1}{m!} \sum_{\substack{\mathbf{Y}_{1}, \dots, \mathbf{Y}_{m} \text{ pos.} \\ \Omega_{1}, \dots, \Omega_{m}}} \varphi((\Omega_{1}, \mathbf{Y}_{1}), \dots, (\Omega_{m}, \mathbf{Y}_{m}))$$

$$= \sum_{k} \frac{1}{k!} \sum_{\substack{\mathbf{X}_{1}, \dots, \mathbf{X}_{k} \text{ pos.} \\ \Omega'_{1}, \dots, \Omega'_{k}}} \varphi_{\mathbf{C}}(\{(\Omega'_{n}, \mathbf{X}_{n}) : \Omega'_{n} \neq \emptyset\}; \{(\emptyset, \mathbf{X}_{n}) : \Omega'_{n} = \emptyset\}) \quad (III.12)$$

$$\times \sum_{l} \frac{1}{l!} \sum_{\substack{\mathbf{Y}_{1}, \dots, \mathbf{Y}_{e} \text{ pos.}}} \varphi((\Omega, \mathbf{Y}_{1}), \dots, (\emptyset, \mathbf{Y}_{l}))$$

with

$$\varphi_{C}((\Omega_{1}, \mathbf{X}_{1}), \dots, (\Omega_{i}, \mathbf{Y}_{i}); (\emptyset, \mathbf{Y}_{1}), \dots, (\emptyset, \mathbf{Y}_{j}))$$

$$= \prod_{1 \leq r < s \leq i} U((\Omega_{r}, \tilde{\mathbf{X}}_{r}), (\Omega_{s}, \tilde{\mathbf{X}}_{s})) \sum_{G_{c}} \prod_{\alpha \in G_{c}} A(\alpha) S_{i+j}((\Omega_{1}, \mathbf{X}_{1}), \dots, (\emptyset, \mathbf{Y}_{j})).$$
(III.13)

The sum \sum_{G_c} is over all graphs whose lines α contain at least one $(\emptyset, \mathbf{Y}_r)$ and in which every $(\emptyset, \mathbf{Y}_r)$ is contained in a connected component containing some (Ω_s, \mathbf{X}_s) . The second factor in (III.12) is now equal to the denominator in (III.3) and thus

$$\langle \Phi \rangle = \sum_{\substack{\{\mathbf{X}_1, \dots, \mathbf{X}_i\} \text{ pos.} \\ \Omega \cap X_n \neq 0, \ \Omega \subset \mathbf{X}_1 \cup \dots \cup \mathbf{X}_i}} \sum_{\substack{\mathbf{Y}_1, \dots, \mathbf{Y}_j \text{ pos.} \\ j!}} \frac{1}{j!}$$

$$\times \varphi_{\mathbf{C}}((\Omega \cap X_1, \mathbf{X}_1), \dots, (\Omega \cap X_i, \mathbf{X}_i); (\emptyset, \mathbf{X}_1), \dots, (\emptyset, \mathbf{Y}_j)).$$
(III.14)

Lemma 9. Given $q \ge 1$ there is an $\varepsilon > 0$ such that for $\lambda > 0$ sufficiently small

and $i, j \ge 1$

$$\sum_{\substack{\mathbf{Y}_{1},\dots,\mathbf{Y}_{j}\\ \sum_{s}|\mathbf{y}_{s}|=l^{2}N}} |\varphi_{C}((\Omega \cap X_{1},\mathbf{X}_{1}),\dots,(\Omega \cap X_{i},\mathbf{X}_{i});(\emptyset,\mathbf{Y}_{1}),\dots,(\emptyset,\mathbf{Y}_{j}))|_{q}$$

$$\leq j! \prod_{r=1}^{i} C(\Omega \cap X_{r},\mathbf{X}_{r}) \exp\left(2l^{-2}\sum_{r=1}^{i}|X_{r}|-\varepsilon lN\right)$$
(III.15)

where $C(\Omega \cap X, \mathbf{X})$ is the bound on $S_{\Omega \cap X, \mathbf{X}}$ obtained in Lemma 8.

Proof. First we sum in (III.13) over graphs G_c with $\{s:((\Omega_1, \tilde{X}_1), (\emptyset, \tilde{Y}_s)) \in G_c\} = I$ fixed. For such graphs vertices $(\emptyset, \mathbf{Y}_s)$, $s \notin I$ must be connected directly or indirectly to $\{(\Omega_r, \mathbf{X}_r): r = 2, ..., i\} \cup \{(\emptyset, \mathbf{Y}_s): s \in I\}$, but there is no restriction on lines $((\emptyset, \tilde{\mathbf{Y}}_s), (\emptyset, \tilde{\mathbf{Y}}_s))$ and $((\Omega_r, \tilde{\mathbf{X}}_r), (\emptyset, \tilde{\mathbf{Y}}_s))$, $s, s' \in I$, r = 2, ..., i. Then by summing over $I \subset \{1, ..., j\}$ we obtain the recursion relation

$$\varphi_{\mathbf{C}}((\Omega_{1}, \mathbf{X}_{1}), \dots, (\Omega_{i}, \mathbf{X}_{i}); (\emptyset, \mathbf{Y}_{1}), \dots, (\emptyset, \mathbf{Y}_{j}))$$

$$= \sum_{\mathbf{I} \subset \{1, \dots, j\}} \prod_{\mathbf{r}=2}^{i} \mathbf{U}((\Omega_{1}, \tilde{\mathbf{X}}_{1}), (\Omega_{\mathbf{r}}, \tilde{\mathbf{X}}_{\mathbf{r}})) \prod_{\mathbf{s} \in \mathbf{I}} \mathbf{A}((\Omega_{1}, \tilde{\mathbf{X}}_{1}), (\emptyset, \tilde{\mathbf{Y}}_{\mathbf{s}}))$$

$$\times (\mathbf{S}_{1} \times \varphi_{\mathbf{C}})((\Omega_{1}, \mathbf{X}_{1}), \dots, (\Omega_{i}, \mathbf{X}_{i}), \{(\emptyset, \mathbf{Y}_{\mathbf{s}})\}_{\mathbf{s} \in \mathbf{I}}; \{(\emptyset, \mathbf{Y}_{\mathbf{s}})\}_{\mathbf{s} \notin \mathbf{I}}).$$
(III.16)

To prove (III.15) we proceed by induction over i+j. By setting $\varphi_{C}(\emptyset; (\emptyset, \mathbf{Y}_{1}), \dots, (\emptyset, \mathbf{Y}_{j})) = 0$ we can start with i+j=1, and clearly $\varphi_{C}((\Omega_{1}, \mathbf{X}_{1}); \emptyset) = S_{\Omega_{1}, \mathbf{X}_{1}}$ satisfies (III.15). Notice that $A((\Omega_{1}, \mathbf{X}_{1}), (\emptyset, \mathbf{Y}_{s})) \neq 0$ only if Y_{s} overlaps or surrounds X_{1} . We say Y_{s} surrounds X_{1} if X_{1} is contained in a hole of Y_{s} and we will write Y_{s} os X_{1} . The number of clusters \mathbf{Y}_{s} satisfying this condition with X_{1} and $|Y_{s}| = l^{2}M$ fixed is bounded by $l^{-2} |X_{1}| \exp O(1)M$.

Define $f(\tilde{\mathbf{X}}) = \sup \{ |f(\mathbf{X})|_q : \mathbf{X} \in \tilde{\mathbf{X}} \}$. Then by using Lemma 8 (take $3\varepsilon \leq 2\delta$) and the induction hypothesis we obtain

$$\begin{split} &\sum_{\substack{\mathbf{Y}_{1},...,\mathbf{Y}_{i} \\ \sum_{s}|\mathbf{Y}_{s}|=l^{2}N}} |\varphi_{C}((\Omega_{1},\mathbf{X}_{1}),\ldots,(\Omega_{i},\mathbf{X}_{i});(\emptyset,\mathbf{Y}_{1}),\ldots,(\emptyset,\mathbf{Y}_{j}))|_{q} \\ &\leq \sum_{I \in \{1,...,j\}} \sum_{\substack{\mathbf{Y}_{1},...,\mathbf{Y}_{i} \\ \sum_{s}|\mathbf{Y}_{s}|=l^{2}N}} 2^{|I|}S_{1}((\Omega_{1},\tilde{\mathbf{X}}_{1})) \\ &\times \varphi_{C}(\{(\Omega_{n},\tilde{\mathbf{X}}_{n})\}_{n=2,...,i},\{(\emptyset,\tilde{\mathbf{Y}}_{s})\}_{s\in I};\{(\emptyset,\mathbf{Y}_{s})\}_{s\notin I}) \\ &\leq S_{1}((\Omega_{1},\tilde{\mathbf{X}}_{1})) \sum_{\substack{\mathbf{Y}_{1},...,\mathbf{Y}_{i} \\ \sum_{s}|\mathbf{Y}_{s}|=l^{2}N}} [\varphi_{C}(\{(\Omega_{n},\tilde{\mathbf{X}}_{n})\}_{n=2,...,i},\{(\emptyset,\mathbf{Y}_{m})\}_{m=1,...,j};\emptyset) \\ &+ 2^{j}\varphi_{C}(\{(\Omega_{n},\tilde{\mathbf{X}}_{n})\}_{n=2,...,i},\{(\emptyset,\mathbf{Y}_{m})\}_{m=1,...,j};\emptyset) \\ &+ \sum_{\substack{I \in \{1,...,j\} \\ 0 < |I| < j}} 2^{|I|}S_{1}((\Omega_{1},\tilde{\mathbf{X}}_{1})) \sum_{\substack{k=1 \\ \sum_{s}|\mathbf{Y}_{s}|=l^{2}k}} \sum_{\substack{(\mathbf{Y}_{s})_{s\notin I} \\ \sum_{s}|\mathbf{Y}_{s}|=l^{2}(N-k)}} \sum_{\substack{(\mathbf{Y}_{s})_{s\notin I} \\ \mathbf{Y}_{s} \text{ os } \mathbf{X}_{1}}} (\varphi_{C}(\{(\Omega_{n},\tilde{\mathbf{X}}_{n})\}_{n=2,...,i},\{(\emptyset,\mathbf{Y}_{s})\}_{s\in I};\{(\emptyset,\mathbf{Y}_{s})\}_{s\notin I}\}) \end{split}$$

Vol. 53, 1980 Particles exist in the low temperature φ_2^4 model

$$\begin{split} &\leq j! \prod_{r=1}^{i} C(\Omega_{r}, \mathbf{X}_{r}) \exp\left(2l^{-2} \sum_{r=2}^{i} |X_{r}|\right) \left[e^{-\varepsilon lN} \right. \\ &+ \sum_{1 < |I|} \frac{1}{|I|!} \sum_{k=1}^{N} \sum_{\substack{\sigma_{s} \geq 1 \\ \sum_{1 \leq s \leq |I|} \sigma_{s} = k}} (2l^{-2} |X_{1}|)^{|I|} e^{(O(1)+2-3\varepsilon l)k} e^{-\varepsilon l(N-k)} \right] \\ &\leq j! \prod_{r=1}^{i} C(\Omega_{r}, \mathbf{X}_{r}) \exp\left(2l^{-2} \sum_{r=2}^{i} |X_{r}| - \varepsilon lN\right) \\ &\times \left[1 + \sum_{1 < |I|} \frac{1}{|I|!} \left((2l^{-2} |X_{1}|)^{|I|} \sum_{\sigma_{1}, \dots, \sigma_{|I|} \geq 1} \exp\left(-2\varepsilon l \sum_{s} \sigma_{s}\right)\right)\right] \\ &\leq j! \prod_{r=1}^{i} C(\Omega_{r}, \mathbf{X}_{r}) \exp\left(2l^{-2} \sum_{r=1}^{i} |X_{r}| - \varepsilon lN\right). \end{split}$$

Proof of Lemma 5. By inserting the bounds (III.15) into (III.14) we obtain

$$\begin{split} |\langle \Phi \rangle|_{q} &\leq \sum_{\substack{\{\mathbf{X}_{1},\dots,\mathbf{X}_{i}\}\\\Omega \cap X_{n} \neq \emptyset, \ \Omega \subset X_{1} \cup \dots \cup X_{i}}} \sum_{\mathbf{Y}_{1},\dots,\mathbf{Y}_{i}} \frac{1}{j!} \left| \varphi_{C}((\Omega \cap X_{1}, \mathbf{X}_{1}), \dots, (\Omega \cap X_{n}, \mathbf{X}_{i}); (\emptyset, \mathbf{Y}_{1}), \dots, (\emptyset, \mathbf{Y}_{j})) \right|_{q} \\ &\leq O(1) \sum_{\substack{\{\mathbf{X}_{1},\dots,\mathbf{X}_{i}\}\\X_{m} \cap X_{n} \neq \emptyset, \ \Omega \subset X_{1} \cup \dots \cup X_{i}}} \prod_{r=1}^{i} C(\Omega \cap X_{r}, \mathbf{X}_{r}) \exp\left(2l^{-2}\sum_{r=1}^{i} |X_{r}|\right) \\ &\leq C_{1}(m_{\Phi}, q)C_{2}(n_{\Phi}, q) \sum_{k=1}^{\infty} \binom{m_{\Phi}(\Omega) + n_{\Phi}(\Omega)}{k} \left(\sum_{\sigma=1}^{\infty} \exp\left(O(1) + 2 - 3\varepsilon l\right)\sigma\right)^{k} \\ &\leq C_{1}(m_{\Phi}, q)C_{2}(n_{\Phi}, q)2^{m_{\Phi}(\Omega) + n_{\Phi}(\Omega)}. \end{split}$$
(III.17)

The third inequality follows since the number of clusters **X** with $|\mathbf{X}| = l^2 \sigma$ and containing a fixed square $\Delta \subset \Omega$ is bounded by $e^{O(1)\sigma}$, while the number of choices of k squares in Ω is bounded by $\binom{m_{\Phi}(\Omega) + n_{\Phi}(\Omega)}{k}$. The factor $2^{m_{\Phi}(\Omega) + n_{\Phi}(\Omega)}$ can be absorbed into C_1C_2 . This proves Lemma 5.

Proof of Lemma 6. We insert (III.14) into (II.6) and obtain the expansion

$$\begin{split} \langle \Phi(\mathbf{R}_{1}); \dots; \Phi(\mathbf{R}_{s}) \rangle &= -\sum_{p=1}^{s} \frac{(-1)^{p}}{p} \sum_{\substack{z_{1} \cup \cdots \cup \sigma_{p} = \{1, \dots, s\} \\ \sigma_{i} \neq \emptyset}} \prod_{k=1}^{p} \\ & \times \sum_{\substack{\{\mathbf{X}_{1}^{k}, \dots, \mathbf{X}_{i_{k}}^{k}\} \text{pos.} \\ [\Omega^{k} \cap X_{j}^{k} = \phi, \Omega^{k} \subset X_{1}^{k} \cup \dots \cup X_{i_{k}}^{k}}} \sum_{\substack{\{\mathbf{Y}_{1}^{k}, \dots, \mathbf{Y}_{j_{k}}^{k}\} \text{ pos.} \\ \times \varphi_{C}^{k}((\Omega^{k} \cap X_{1}^{k}, \mathbf{X}_{1}^{k}), \dots, (\Omega^{k} \cap X_{i_{k}}^{k}, \mathbf{X}_{i_{k}}^{k}); (\emptyset, \mathbf{Y}_{1}^{k}), \dots, (\emptyset, \mathbf{Y}_{j_{k}}^{k}))] \end{split}$$

where for the definition of φ_C^k the fields $\Phi_{\Omega'}$ in (II.12) are replaced by $\Phi(\bigcup_{i \in \sigma_k} R_k)_{\Omega'_k}$. By the proofs of the Lemmas 5 and 9 we know that this sum is absolutely convergent, also if the φ_C^k are expressed as a sum of products of functions S_1 . We may thus sum up first all contributions from sets $C = (\mathbf{X}_1^1, \ldots, \mathbf{X}_{i_p}^p, \mathbf{Y}_1^1, \ldots, \mathbf{Y}_{j_p}^p)$ which factorize, i.e. for which we have two regions

 $G_1, G_2 \subseteq \mathbb{R}^2$ without hole separated by a positive distance, and a partition (C_1, C_2) of C such that $Z \subseteq G_1$ for $Z \in C_1$ and $Z \subseteq G_2$ for $Z \in C_2$. Notice that if the argument A of a φ_C^k only contains clusters of a factorizing set C, then $\varphi_C^k(A)$ factorizes accordingly to C into a product $\varphi_C^k(A_1) \cdot \varphi_C^k(A_2)$. From a standard argument using formal power series [R], [EMS] we conclude that this first sum is zero. From the remaining terms we can extract a factor $e^{-\delta d(\Omega_1,...,\Omega_s)}$ by using Lemma 8. The assertion now follows since $C_1(\cdot, q), C_2(\cdot, q)$ are concave and since $|f_1 \otimes f_2|_q = |f_1|_q \cdot |f_2|_q$ and

$$\sum_{\sigma_1\cup\cdots\cup\sigma_p=\{1,\ldots,s\}}\frac{1}{p}\leq s^s.$$

IV. Proof of Lemma 7

We assume that $|\Sigma \cap N(\mathfrak{L}_i)| = 0$ for all $i \in I$ (otherwise $F_{\Omega', \mathbf{X}, \Sigma} = 0$) and omit the r-factors since they are bounded by

$$\left| \prod_{i \in \mathfrak{Y}} r_i^{|\Sigma \cap N(\mathfrak{L}_i)|} \right| \leq ((e^{3L})^{(2L+1)l^{-1}})^{|\Sigma|}$$
$$\leq e^{7l^{15}|\Sigma|} e^{o(\lambda^{-1/2})|\Sigma|}$$

and will be dominated by a factor $e^{-b\lambda^{-1/2}|N(\Sigma)|}$, b > 0. Let $B_0 = \{(j, \alpha) : j \in l\mathbb{Z}^4 \cap X, \alpha \in \pi_0, \alpha(i) = 0 \text{ if } \mathfrak{Q}_i \cap X = \emptyset\}$. By expanding the product over $\beta \in B_0$ we obtain the following expression for $\delta_s^{\Gamma} F_{\Omega', \mathbf{X}, \Sigma}(\mathbf{r}, \mathbf{t}, \mathbf{h}, \mathbf{s})$

$$\begin{split} \delta_{s}^{\Gamma} F_{\Omega', \mathbf{X}, \Sigma} \\ &= \sum_{\mathbf{B} \subset \mathbf{B}_{0}} \delta_{s}^{\Gamma} \int \! d\psi_{C(t,s)} e^{-F(\Lambda \cap \mathbf{X}, \Sigma)} \prod_{\beta \in \mathbf{B}} \left[h(\alpha) \partial_{t}^{\alpha} C_{j}(t,s) \cdot \Delta_{\phi} \right] \! \Phi_{\Omega'} \chi_{\Sigma} e^{-V(\Lambda \cap \mathbf{X})}. \end{split}$$

The derivatives are computed as in (II.1), leading to

$$\begin{split} \delta_{s}^{\Gamma} F_{\Omega',\mathbf{X},\Sigma} &= \sum_{\Gamma_{1} \cup \Gamma_{2} = \Gamma} \sum_{B \subset B_{0}} \int_{\{0\}}^{\{1\}} ds \int d\psi_{C(t,s(\Gamma))} \bigg[\sum_{\pi \in P(\Gamma_{2})} \prod_{\gamma \in \pi} \partial_{s}^{\gamma} C(t,s(\Gamma)) \cdot \Delta_{\psi} \bigg] \\ &\times e^{-F(\Lambda \cap \mathbf{X},\Sigma)} \bigg[\sum_{\bigcup_{\beta \in B} \gamma_{\beta} = \Gamma_{2}} \prod_{\beta \in B} h(\alpha) \partial_{s}^{\gamma_{\beta}} \partial_{t}^{\alpha} C_{j}(t,s(\Gamma)) \cdot \Delta_{\phi} \bigg] \Phi_{\Omega'} \chi_{\Sigma} e^{-V(\Lambda \cap \mathbf{X})} \end{split}$$

where $P(\Gamma_2)$ is the set of partitions of Γ_2 . We also write the kernels $\partial_s^{\gamma}C$ as a sum over localizations in the $l\mathbb{Z}^2$ lattice

$$\begin{split} \delta_{s}^{\Gamma} F_{\Omega', \mathbf{X}, \Sigma} &= \Sigma' \int d\psi_{C(t, s(\Gamma))} \bigg[\prod_{\gamma \in \pi} \partial_{s}^{\gamma} C_{j, \gamma}(t, s(\Gamma)) \cdot \Delta \psi \bigg] e^{-F(\Lambda \cap \mathbf{X}, \Sigma)} \\ & \times \bigg[\prod_{\beta \in \mathbf{B}} h(\alpha) \partial_{s}^{\gamma_{B}} \partial_{t}^{\alpha} C_{j}(t, s(\Gamma)) \cdot \Delta_{\phi} \bigg] \Phi_{\Omega'} \chi_{\Sigma} e^{-\mathbf{V}(\Lambda \cap \mathbf{X})} \end{split}$$
(IV.1)

where

$$\Sigma' = \sum_{\Gamma_1 \cup \Gamma_2 = \Gamma} \sum_{\mathbf{B} \subset \mathbf{B}_0} \sum_{\bigcup_{\beta \in \mathbf{B}} \gamma_\beta = \Gamma_1} \sum_{\pi \in \mathbf{P}(\Gamma_2)} \sum_{\{j_\gamma\}_{\gamma \in \pi}}.$$

Each term in this sum is of the form

$$\Sigma'' \int d\psi_{C(t,s(\Gamma))} W \Phi' \chi'_{\Sigma} e^{-V(\Lambda \cap X) - F(\Lambda \cap X,\Sigma)}.$$
 (IV.2)

 Σ'' is obtained and interpreted as follows. By definition $d/d\phi: \phi^n:=n:\phi^{n-1}$: and this is rewritten as $\Sigma: \phi^{n-1}$:. This convention is inductively used for all derivatives $d/d\phi$ acting on $\Phi_{\Omega'}\chi_{\Sigma}e^{-\nu}$ or on derivatives of this expression. Every field ϕ is then rewritten as a function of ψ , $\phi(x) = \psi(x) + (\xi - g(x))$, producing again a sum of terms. Now we perform the derivatives $d/d\psi$ in the same way.

In order to bound the norm $|\delta_s^{\Gamma} F_{\Omega', \mathbf{X}, \Sigma}|_q$, q > 1, associated to a given partition π_1 of $\{1, \ldots, N\}$ we bound the product

$$\delta_{s}^{\Gamma} F_{\Omega', \mathbf{X}, \Sigma}[\omega] = \langle (\delta_{s}^{\Gamma} F_{\Omega', \mathbf{X}, \Sigma})_{\pi_{1}}, \omega \rangle$$

for functions $\omega \in L^p$, $p^{-1} + q^{-1} = 1$.

By Hölder's inequality

$$\left\|\delta_{s}^{1}F_{\Omega',\mathbf{X},\Sigma}[\omega]\right\| \leq \Sigma'\Sigma'' \left\|W\Phi'[\omega]\right\|_{r} \left\|\chi'_{\Sigma}\exp\left(-V(\Lambda\cap X) - F(\Lambda\cap X,\Sigma)\right)\right\|_{r'} \qquad (IV.3)$$

where r and r' are dual Hölder indices and r is some sufficiently large even integer $(r \approx 35\ 000)$. The last factor in (IV.3) is bounded as in [GJS III] by

$$\begin{aligned} \|\chi'_{\Sigma} e^{-V(\Lambda \cap X) - F(\Lambda \cap X, \Sigma)}\|_{r'} \\ \leq \left(\prod_{\Box \in X \cap \Lambda} \nu(\Box)!\right) \exp\left(a'\lambda^{1/2} |X \cap \Lambda|\right) \exp\left(-3b\lambda^{-1/2}(|N(\Sigma)| + |X'|)\right) \\ (IV.4) \end{aligned}$$

for some positive constants a', b, where $\nu(\Box)$ is the number of times $\chi_{\Sigma(\Box)}$ has been derived, and where

$$\mathbf{X}' = \bigcup_{\boldsymbol{\nu}(\Box) > 0} \Box.$$

Each covariance is now written as a sum of terms localized (in both variables) in unit lattice squares. This induces an expansion of $W\Phi'[\omega]$ into a sum

$$W\Phi'[\omega] = \sum_{u} W_{u}\Phi'[\omega]$$
 (IV.5)

of Wick monomials localized in unit squares. The resulting number of terms is bounded by

$$\sum_{u} 1 \le \prod_{\gamma \in \pi} l^4 \prod_{\beta \in B} l^4$$
(IV.6)

and the monomials in W_u are smeared out with localized products of functions $(\xi - g)$ and $\partial_s^{\gamma} \partial_t^{\alpha} C$.

In order to bound the resulting graphs and the sum Σ' we use the decay properties of kernels $\partial_s^{\gamma} \partial_t^{\alpha} C_j$ established in [S II], adapted to the case of unit mass and large length scale. Define

$$\begin{split} &\mathbf{i}(\alpha) = \min \{\mathbf{i} \in l\mathbb{Z} : \alpha(\mathbf{i}) \neq 0\}, \\ &\mathbf{d}(\beta) = \mathbf{d}(\mathbf{j}, \alpha) = \operatorname{dist} \left(\Delta_{\mathbf{j}_1}, \mathfrak{L}_{\mathbf{i}(\alpha)}\right) + \operatorname{dist} \left(\Delta_{\mathbf{j}_2}, \mathfrak{L}_{\mathbf{i}(\alpha)}\right) + \operatorname{dist} \left(\Delta_{\mathbf{j}_1}, \Delta_{\mathbf{j}_2}\right), \end{split}$$

 $d(j, \gamma) = \max \{ \text{dist} (\Delta_{j_1}, b) + \text{dist} (\Delta_{j_2}, b) : b \in \gamma \},\$ $d(\gamma, \alpha) = \min \{ \text{dist} (b, \mathfrak{L}_{i(\alpha)}) : b \in \gamma \},\$ $d(\gamma) : \text{the length of the shortest path in } \mathbb{R}^2 \text{ connecting all } b \in \gamma.$

Lemma 10. Let $p \ge 1$ and $\varepsilon > 0$ be given. Then for l_0 sufficiently large there are positive constants $M(\gamma, \beta)$, M_3 and c such that for $l \ge l_0$.

$$\begin{aligned} \|\partial_{s}^{\gamma}\partial_{t}^{\alpha}C_{j}(t,s)\|_{p} &\leq l^{4/p}M_{2}(\gamma,\beta)e^{-(1-\varepsilon)d(\alpha)} \\ &\times e^{-6cd(j,\alpha)}e^{-5c(2d(\gamma)+3d(j,\gamma))} \end{aligned} \tag{IV.7}$$

and (see also [GJS I])

$$\sum_{\boldsymbol{\gamma}\in\mathbf{P}(\Gamma_2)}\prod_{\boldsymbol{\gamma}\in\boldsymbol{\pi}}\mathbf{M}_2(\boldsymbol{\gamma},(\alpha=0,j)) \leq e^{\mathbf{M}_3|\Gamma_2|l^{-1}},$$
(IV.8)

$$\sum_{\bigcup_{\beta \in \mathbf{B}} \gamma_{\beta} = \Gamma_{1}} M_{2}(\gamma_{\beta}, \beta) \leq e^{M_{3}|\mathbf{X}|l^{-2}}.$$
 (IV.9)

Furthermore one has

π

$$\sum_{\{\mathbf{j}\gamma\}_{\gamma\in\pi}}\prod_{\gamma\in\pi}e^{-\mathrm{cd}(\mathbf{j}_{\gamma},\gamma)} \leq e^{\mathbf{M}_{\gamma}|\Gamma_{2}|l^{-1}},\tag{IV.10}$$

$$\sum_{\mathbf{B}\subset\mathbf{B}_{0}}\prod_{\boldsymbol{\beta}\in\mathbf{B}}e^{-\mathbf{cd}(\mathbf{j},\boldsymbol{\alpha})} \leq e^{\mathbf{M}_{0}|\mathbf{X}|l^{-2}}$$
(IV.11)

for some constants $M_0, M_7 > 0$. (IV.7) remains true if the kernel is regarded as a function of a single variable $\partial_s^{\gamma} \partial_t^{\alpha} C_j(t, s)(x, x), \ \gamma \neq 0$ or $\alpha \neq 0$.

By using Lemma 10 we bound $||W\Phi'[\omega]||_r$ times the first factor of (IV.4) as follows. Let

$$\begin{split} \mathbf{P}(\Delta) &= \{ (\mathbf{j}, \alpha) \in \mathbf{B} : \Delta \subset \Delta_{\mathbf{j}_1} \cup \Delta_{\mathbf{j}_2} \} \cup \{ (\mathbf{j}_{\gamma}, \gamma) : \gamma \in \pi, \ \Delta \subset \Delta_{\mathbf{j}_{\gamma,1}} \cup \Delta_{\mathbf{j}_{\gamma,2}} \} \\ \mathbf{M}(\Delta) &= \operatorname{card} \mathbf{P}(\Delta), \\ \mathbf{M}_{-}(\Delta) &= \operatorname{card} \{ (\mathbf{j}, \alpha) \in \mathbf{B} : \Delta \subset \Delta_{\mathbf{j}_1} \cup \Delta_{\mathbf{j}_2}, \ (\xi - g) \upharpoonright \Delta \neq 0 \}, \\ \mathbf{M}_{+}(\Delta) &= \mathbf{M}(\Delta) + \mathbf{m}_{\Phi_{\alpha}}(\Delta), \end{split}$$

and let $m = m_{\Phi_{\Omega'}}$, $n = n_{\Phi_{\Omega'}}$. Then

$$\begin{split} &\prod_{\square=\mathbf{X}\cap\Lambda} \nu(\square)! \| W\Phi'[\boldsymbol{\omega}] \|_{r} \leq \| \boldsymbol{\omega} \|_{p} C_{3}(0, \Omega', q) e^{26l^{4} |\mathbf{X}'|} \\ &\times C(\mathbf{X}, \Sigma)^{n(\Omega')} e^{(1/2+6c)ln(\Omega')} \prod_{\square=\mathbf{X}\cap\Lambda} \nu(\square)! [(rn(\square))!]^{1/r} \\ &\times \prod_{\Delta \in \mathbf{X}} \left[(4r(\mathbf{M}_{+}(\Delta) - \sum_{\square=\Delta} \nu(\square)))! \right]^{1/r} \mathbf{M}_{1}(\mathbf{p}, \mathbf{r})^{\mathbf{M}_{+}(\Delta)} C(\mathbf{X}, \Sigma)^{3\mathbf{M}_{-}(\Delta)} \\ &\times \prod_{\gamma \in \pi} \mathbf{M}_{2}(\gamma, (0, j_{\gamma})) e^{-5cd(j_{\gamma}, \gamma)} e^{-2cl |\gamma|} \\ &\times \prod_{\beta \in \mathbf{B}} \mathbf{M}_{2}(\gamma_{\beta}, \beta) e^{-5cd(\beta)} e^{-2cl |\gamma_{\beta}|} \end{split}$$
(IV.12)

where we have made the substitution

$$h(\alpha)e^{-(1-\varepsilon)d(\alpha)}e^{-cd(j,\alpha)}$$

$$\times \prod_{\gamma \in \pi} l^{8}e^{-10[d(\gamma)+d(j_{\gamma},\gamma)]} \prod_{\beta \in B} l^{8}e^{-10c[d(\gamma_{\beta})+d(j,\gamma_{\beta})]}$$

$$\rightarrow e^{26l^{4}|X'|}e^{(1/2+6c)\ln(\Omega')}e^{-2cl|\Gamma|}C_{3}(0,\Omega',q).$$
(IV.13)

Except for this substitution and for the factors $C(\mathbf{X}, \Sigma)$ arising when $\phi \neq \psi$, (IV.13) is a standard estimate (see [GJS I], but with separate copies of unit squares for the monomials of W_u and $\Phi'[\omega]$). The convergence factors allowing for (IV.13) are obtained as follows. We use

(1) for terms from contractions $\partial_s^{\gamma} C_{j_{\gamma}} \cdot \Delta_{\psi}$

$$l^{8} \mathrm{e}^{-10\mathrm{c}[\mathrm{d}(\gamma)+\mathrm{d}(\mathrm{j}_{\gamma},\gamma)]} \leq \mathrm{e}^{-2\mathrm{c}l\,|\gamma|} \quad \text{if} \quad \mathrm{d}(\gamma)+\mathrm{d}(\mathrm{j}_{\gamma},\gamma) > 0.$$

If $d(\gamma) + d(j_{\gamma}, \gamma) = 0$ then $|\gamma| \le 4$ and $F(\Delta_{j_{\gamma,1}} \cup \Delta_{j_{\gamma,2}}, \Sigma) = 0$, and we use

 $\lambda^{\epsilon/4} l^4 < e^{-cl |\gamma|}$ when the contraction is to V.

 $l^4 < e^{-cl |\gamma|} e^{l^4 |\Box \cap X'|}$ when the contraction is to $\chi_{\Sigma(\Box)}$.

 $l^4 < e^{-cl |\gamma|} e^{5cl}$ when the contraction is to $\Phi_{\Omega'}$.

The same holds if (j_{γ}, γ) is replaced by (j, γ_{β}) . (2) for terms from contractions $\partial_t^{\alpha} C_i \cdot \Delta_{\phi}$

$$\begin{aligned} &(|\mathbf{h}(\alpha)| \le \mathrm{e}^{(1-2\varepsilon)(\mathrm{d}(\alpha)+l)}) \\ &l^{8}\mathrm{e}^{-3\varepsilon l} < 1 \quad \text{if} \quad \mathrm{d}(\alpha) \ge \varepsilon^{-1}l. \\ &l^{8}\mathrm{e}^{l}\mathrm{e}^{-\mathrm{cd}(\mathbf{j},\alpha)} < 1 \quad \text{if} \quad \mathrm{d}(\mathbf{j},\alpha) \ge 2\mathrm{c}^{-1}l. \end{aligned}$$

If $d(\alpha) < \varepsilon^{-1}l$ and $d(j, \alpha) < 2c^{-1}l$ then we use

 $\lambda^{\epsilon/4} l^4 e^{l/2} < 1$ when the contraction is to V.

 $l^4 e^{l/2} < e^{l(1/2+c)}$ when the contraction is to $\Phi_{\Omega'}$.

 $l^4 e^{l/2} < (e^{l^4 |\Box \cap X'|})^{1/2Kl^2}$ for $K = 24/c\varepsilon$, when the contraction is to $\chi_{\Sigma(\Box)}$. This is sufficient since the number of β 's in B such that $d(\alpha) < \varepsilon^{-1}l$ and $d(j, \alpha) < 2c^{-1}l$ and $d(j, \alpha) < 2c^{-1}l$ and $\Box \subset \Delta_{j_1} \cup \Delta_{j_2}$ is bounded by Kl^2 .

(3) If a contraction $d/d\phi$ from $\Phi_{\Omega'}$ is to V then we get a factor $\lambda^{1/2-\epsilon}$, and if it is to $\chi_{\Sigma(\Box)}$ we use that

 $1 \leq \lambda^{1/2-\varepsilon} e^{l^4 |\Box \cap X'|}$

Next we use a similar argument to cancel the factor $C(\mathbf{X}, \Sigma)^{3M_{-}(\Delta)}$. Namely for each contraction $\partial_s^{\gamma_{\beta}} \partial_t^{\alpha} C_j \cdot \Delta_{\phi}$ with $(\xi - g) \upharpoonright (\Delta_{j_1} \cup \Delta_{j_2}) \neq 0$ we have $d(\beta) > L$ and thus

 $C(\mathbf{X}, \Sigma)^6 e^{-cd(\beta)} < 1$

since $C(\mathbf{X}, \Sigma) = O(\lambda^{-1/2})$ and $L \simeq (\log \lambda)^2$.

We continue by estimating (IV.3). The sum Σ' is controlled by (IV.8),..., (IV.11) and a factor $e^{-cl |\Gamma|}$ (notice that $\sum_{\Gamma_1 \cup \Gamma_2 = \Gamma} 1 = 2^{|\Gamma|}$). There remains a factor

$$e^{-cl |\Gamma|} e^{-b\lambda^{-1/2}(|\mathbf{N}(\Sigma)| + |\mathbf{X}'|)} e^{a''l^{-2} |\mathbf{X} \cap \Lambda|} \|\boldsymbol{\omega}\|_{\mathbf{p}} C_3(0, \Omega', \mathbf{q})$$
(IV.14)

times the maximum over admissible Γ_i , B, γ_β , π and j_γ of

$$\Sigma'' \prod_{\Box} \nu(\Box)! \prod_{\Delta} \left[\left(4r \left(M_{+}(\Delta) - \sum_{\Box \subset \Delta} \nu(\Box) \right) \right)! \right]^{1/r} M_{1}^{M_{+}(\Delta)} \times \prod_{\beta \in B} e^{-4cd(\beta)} \prod_{\gamma \in \pi} e^{-4cd(j_{\gamma}, \gamma)},$$
(IV.15)

where we have put the factors

$$C(\mathbf{X}, \Sigma)^{n(\Omega')} e^{(1/2+6c)\ln(\Omega')} \prod_{\Box} [(rn(\Box))!]^{1/r} \cdot C_3(0, \Omega', q)$$

into $C_3(n, \Omega', q)$ by using $(ab)! \le a^{ab}(b!)^a$. This inequality together with the facts that $a! b! \le (a+b)!$, $a! \le a^a$ and $(a+b)! \le (2a)! (2b)!$ implies that

$$\prod_{\Box} \nu(\Box)! \prod_{\Delta} \left[\left(4r \left(M_{+}(\Delta) - \sum_{\Box = \Delta} \nu(\Box) \right) \right)! \right]^{1/r} M_{1}^{M_{+}(\Delta)} \\ \leq \prod_{\Delta} (4^{4} r^{4} M_{1})^{M(\Delta) + m(\Delta)} ((2M(\Delta))!)^{4} ((2m(\Delta))!)^{4} \\ \leq e^{O(1)m(\Omega')^{4/3}} \prod_{\Delta} e^{O(1)M(\Delta)^{4/3}}$$
(IV.16)

In a similar way we bound the number of terms coming from differentiations and from expressing V' and $\Phi_{\Omega'}$ in terms of ψ :

$$\Sigma'' 1 \leq \prod_{\Box} 2^{n(\Box)} \prod_{\Delta} K^{M(\Delta)}$$
(IV.17)
 $\times (n(\Delta) + |\Delta \cap X'| + 4)(n(\Delta) + |\Delta \cap X'| + 8) \cdots (n(\Delta) + |\Delta \cap X'| + 4M_{+}(\Delta)).$

By using

$$(N+4)(N+8)\cdots(N+4M) \leq (N+4M)^{M}$$
$$\leq (4M)^{M} \left(1 + \frac{N}{4M}\right)^{M}$$
$$\leq (4M)^{M} e^{N/4}$$

the right side of (IV.17) can be bounded as follows.

$$\Sigma'' 1 \le e^{n(\Omega') + 1/4|X'|} e^{O(1)m(\Omega')^{4/3}} \prod_{\Delta} e^{O(1)M(\Delta)^{4/3}}$$
(IV.18)

Now by putting together $(IV.14), \ldots, (IV.18)$ we arrive at the bound

$$\begin{split} \left| \delta_{s}^{\Gamma} F_{\Omega', \mathbf{X}, \Sigma} [\boldsymbol{\omega}] \right| &\leq \left\| \boldsymbol{\omega} \right\|_{p} C_{1}(m, q) C_{3}(n, \Omega', q) \\ &\times \exp\left(-cl \left| \Gamma \right| - b\lambda^{-1/2} \left| N(\Sigma) \right| + a'' l^{-2} \left| X \cap \Lambda \right| \right) \\ &\times \sup_{\mathbf{B}, \{\gamma_{\beta}\}, \{\gamma\}, \{\mathbf{i}_{\gamma}\}} \prod_{\Delta} \left[e^{\rho M(\Delta)^{4/3}} \prod_{\mathbf{p} \in \mathbf{P}(\Delta)} e^{-2cd(\mathbf{p})} \right] \end{split}$$
(IV.19)

The product $\prod_{\Delta} [\cdot \cdot \cdot]$ in (IV.19) is bounded by using the 'pushout principle' as in [S II], cf. also [GJS I]. This goes as follows. There are not more than $8(a+3)^2$

 $l^{-1} d(\mathbf{p}_k) > \frac{1}{6} k^{1/3} - 3.$

Or, by summing over k

$$l^{-1}\sum_{\mathbf{p}\in\mathbf{P}(\Delta)}d(\mathbf{p}) > \frac{1}{8}M(\Delta)^{4/3} - 3M(\Delta)$$

and thus

$$\rho \mathbf{M}(\Delta)^{4/3} \leq O(1) + \sum_{\mathbf{p} \in \mathbf{P}(\Delta)} \operatorname{cd}(\mathbf{p})$$

This proves Lemma 7.

Acknowledgement

It is a great pleasure for the author to thank Professor J.-P. Eckmann for many discussions and valuable criticism.

REFERENCES

- [BG] T. BALABAN and K. GAWEDZKI, A low temperature expansion for the pseudoscalar Yukawa model of quantum fields in two space-time dimensions. Harvard preprint 1980.
- [Br] D. BRYDGES, A rigorous approach to Debye screening in dilute classical Coulomb systems. Commun. Math. Phys. 58, 313 (1978).
- [BF] D. BRYDES and P. FEDERBUSH, Debye screening. To appear in Commun. Math. Phys.
- [BI] C. BURNAP, Harvard University Thesis 1976.
- [B II] C. BURNAP, Isolated one particle states in boson quantum field theory models. Ann. Phys. 104, 184 (1977).
- [DE I] J. DIMOCK and J.-P. ECKMANN, On the bound state in weakly coupled $\lambda (\phi^6 \phi^4)_2$. Commun. Math. Phys. 51, 41 (1976).
- [DE II] J. DIMOCK and J.-P. ECKMANN, Spectral properties and bound state scattering for weakly coupled $\lambda \mathfrak{P}(\phi)_2$ models. Ann. of Phys., Vol. 103, No 2, 289 (1977).
- [DG] J. DIMOCK and J. GLIMM, Measures on Schwartz distribution space and applications to $\mathfrak{P}(\phi)_2$ field theories. Adv. Math. 12, 58 (1974).
- [EMS] J.-P. ECKMANN, J. MAGNEN and R. SÉNÉOR, Decay properties and Borel summability for the Schwinger functions in $\mathfrak{P}(\phi)_2$ theories. Commun. Math. Phys. 39, 251 (1975).
- [GJ] J. GLIMM and A. JAFFE, The resummation of one particle lines, Commun. Math. Phys. 67, 267 (1979).
- [GJS I] J. GLIMM, A. JAFFE and T. SPENCER, The Wightman axioms and particle structure in the $\mathfrak{P}(\phi)_2$ quantum field model, Ann. Math. 100, 585 (1974).
- [GJS II] J. GLIMM, A. JAFFE and T. SPENCER, The particle structure of the weakly coupled $\mathfrak{P}(\phi)_2$ model and other applications of high temperature expansions. In: G. Velo and A. Wightman Eds, Constructive quantum field theory, Springer Lecture Notes in Physics, Vol. 25, Berlin – Heidelberg – New York (1973).
- [GJS III] J. GLIMM, A. JAFFE and T. SPENCER, A convergent expansion about mean field theory. Ann. Phys. 101, 610 (1976).
- [I] J. IMBRIE, Mass spectrum of the two dimensional $\lambda \varphi^4 \frac{1}{4} \varphi^2 \mu \varphi$ quantum field model. Harvard Preprint 1980.
- [K] H. KOCH, Irreducible kernels and bound states in $\lambda \mathfrak{P}(\varphi)_2$ models. Ann. Inst. H. Poincaré 31, 173 (1979).
- [OS] K. OSTERWALDER and R. SÉNÉOR, The scattering matrix is non-trivial for weakly coupled $\Re(\phi)_2$ models. Helv. Phys. Acta 49, 525 (1976).

- [OSch] K. OSTERWALDER and R. SCHRADER, Axioms for Euclidean Green's functions. Commun. Math. Phys. 31, 83 (1973), 42, 281 (1975).
- [R] D. RUELLE, Statistical mechanics. New York, Benjamin 1969.
- [S I] T. SPENCER, The mass gap for the $\mathfrak{P}(\phi)_2$ field model with a strong external field. Commun. Math. Phys. 39, 63 (1974).
- [S II] T. SPENCER, The decay of the Bethe-Salpeter kernel in $\mathfrak{P}(\phi)_2$ quantum field models. Commun. Math. Phys. 44, 143 (1975).
- [SZ] T. SPENCER and F. ZIRILLI, Scattering states and bound states in $\lambda \mathfrak{P}(\phi)_2$. Commun. Math. Phys. 49, 1 (1976).
- [Su I] S. SUMMERS, Harvard University Thesis 1979.
- [Su II] S. SUMMERS, On the phase diagram of a $\mathfrak{P}(\varphi)_2$ quantum field model. CNRS Preprint 1979.
- [Su III] S. SUMMERS, The decay of the Bethe-Salpeter kernel and isolated one-particle states in strongly coupled φ_2^4 . CNRS Preprint 1980.